

Complex control variates

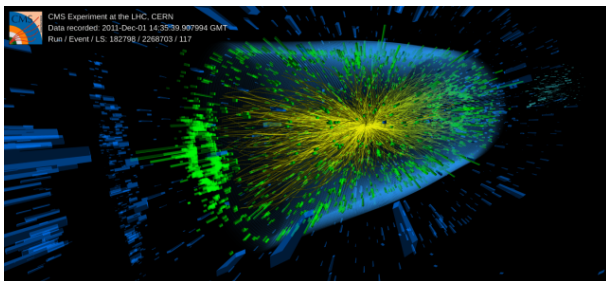
Yukari Yamauchi

- arXiv:2205.12303 [hep-lat] with Scott Lawrence and Hyunwoo Oh
- arXiv:2212.14606 [hep-lat] with Scott Lawrence
- arXiv:2311.13002 [hep-lat] with Scott Lawrence

November 24th, 2023, Large-scale lattice QCD simulation and application of machine learning, university of Tsukuba



Sign problems in lattice QCD



Tom McCauley/CMS/CERN

Path integral

$$\langle \mathcal{O}(t_0) \rangle = \text{Tr} \left[e^{-\beta H} e^{iHt_0} \mathcal{O} e^{-iHt_0} \right] = \frac{1}{Z} \int \mathcal{D}[\psi, U] e^{-S_{\text{SK}}} \mathcal{O}(t_0)$$

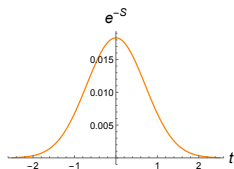
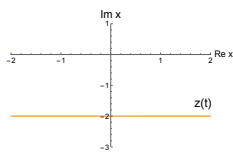
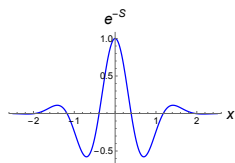
The average phase

$$\langle \sigma \rangle = \frac{\int \mathcal{D}[\psi, U] e^{-S_{\text{SK}}}}{\int \mathcal{D}[\psi, U] |e^{-S_{\text{SK}}}|} \propto e^{-V}$$

Methods

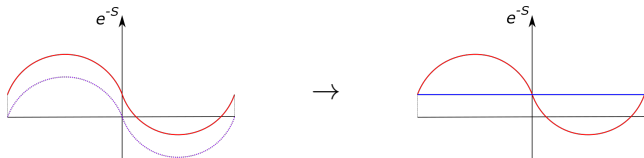
1. Contour deformation / complex normalizing flow

Example: $S(x) = x^2 - 4ix$



Demonstration: scalar fields theories in $0 + 1$ and $1 + 1-d$

2. Complex control variates (subtraction)



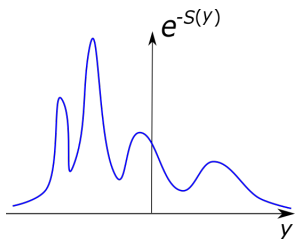
Demonstration: Spin system, Thirring model in $1 + 1-d$

Normalizing flows — when no sign problem

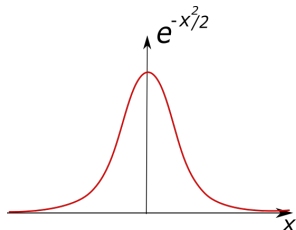
Normalizing flow¹ $\vec{y} = f(\vec{x})$:

$$\det \left(\frac{\partial \vec{y}}{\partial \vec{x}} \right) e^{-S(\vec{y})} = \mathcal{N} \prod_{i=1}^N \frac{1}{\sqrt{2\pi}} e^{-x_i^2/2} = \mathcal{N} G_N(\vec{x})$$

Some distribution



Simple Gaussian distribution



Map
 \leftrightarrow
 $y(x)$

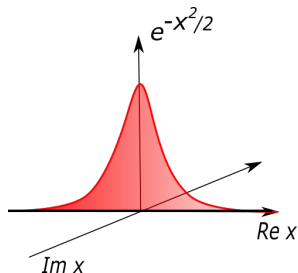
if such a map exists, then

$$\langle \mathcal{O} \rangle = \frac{\int_{\mathbb{R}^N} \mathcal{D}\vec{y} e^{-S(\vec{y})} \mathcal{O}(\vec{y})}{\int_{\mathbb{R}^N} \mathcal{D}\vec{y} e^{-S(\vec{y})}} = \frac{\int_{\mathbb{R}^N} d\vec{x} G_N(\vec{x}) \mathcal{O}(\vec{y}(\vec{x}))}{\int_{\mathbb{R}^N} d\vec{x} G_N(\vec{x})}$$

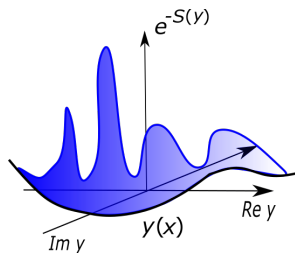
¹M. Albergo et al. Phys. Rev. D 100, 034515(2019)
K. A. Nicoli, et al. Phys. Rev. E 101, 023304(2020)

Complex normalizing flow²/contour deformations

$$\mathcal{N}G_N(\vec{x}) = \det\left(\frac{\partial \vec{y}}{\partial \vec{x}}\right) e^{-S(\vec{y}(\vec{x}))}$$



Map
 \leftrightarrow
 $y(x)$



Expectation values:

$$\frac{\int_{\mathbb{R}^N} d\vec{x} G_N(\vec{x}) \mathcal{O}(y(x))}{\int_{\mathbb{R}^N} d\vec{x} G_N(\vec{x})} = \frac{\int_{\vec{y}(\mathbb{R}^N)} d\vec{y} e^{-S(\vec{y})} \mathcal{O}(\vec{y})}{\int_{\vec{y}(\mathbb{R}^N)} d\vec{y} e^{-S(\vec{y})}} \stackrel{?}{=} \langle \mathcal{O} \rangle$$

Perfect complex normalizing flow \rightarrow No sign problems!

²S. Lawrence and YY, arXiv:2101.05755 [hep-lat]

Machine-learn a map

Contour deformation

Loss function: $L = -\log\langle\sigma\rangle$

The gradient of $\langle\sigma\rangle$ is **sign-free**³!

$$\partial_{\mathbf{v}}(-\log\langle\sigma\rangle) = -\frac{\int_{\mathcal{C}} \mathcal{D}\phi (\partial_{\mathbf{v}} \operatorname{Re} S) |e^{-S}|}{\int_{\mathcal{C}} \mathcal{D}\phi |e^{-S}|}$$

- Needs MCMC sampling
- Can find “good enough” contours

Complex normalizing flow

Loss function:

$$L = \left\langle \left| \det \left(\frac{\partial \mathbf{y}(x_i)}{\partial x_i} \right) e^{-S(\mathbf{y}(x_i))} - \mathcal{N} G_N(x_i) \right| \right\rangle_{G_N}$$

- No MCMC sampling
- Find only very good contours

³A. Alexandru, P. Bedaque, H. Lamm, and S. Lawrence, arXiv:1804.00697 [hep-lat]

Normalizing flows for complex coupling model in $0 + 1-d^4$

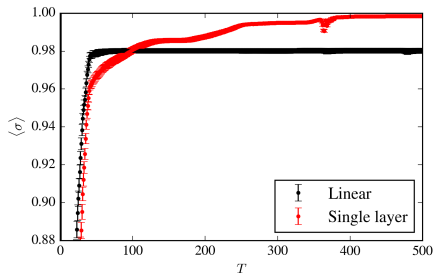
$$S = \sum_{i=1}^N \frac{m}{2} \phi_i^2 + \frac{(\phi_i - \phi_{i-1})^2}{2} + \lambda \phi_i^4, \text{ with } \lambda \in \mathbb{C}$$

Neural network $\vec{\phi}(\vec{x}) = f(\vec{x}) + ig(\vec{x})$, $\sigma = \text{sigmoid}$

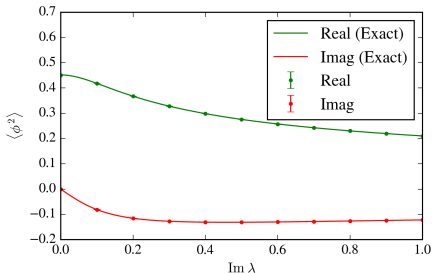
$$f, g(\vec{x}) = (L_{N,N} + L_{N,2N} \otimes \sigma \otimes L_{2N,2N} \cdots L_{2N,2N} \otimes \sigma \otimes L_{2N,N}) \vec{x}$$

Demonstration: $N = 10, m = 0.5$

$\lambda = 0.1 + i$



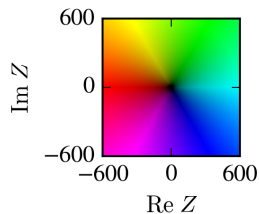
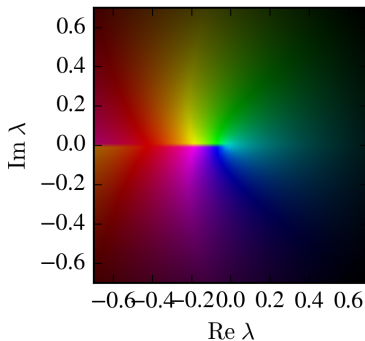
$\text{Re } \lambda = 0.1$



⁴S. Lawrence, H. Oh, and YY, arXiv:2205.12303 [hep-lat]

Partition function in 0 + 1-dimensions

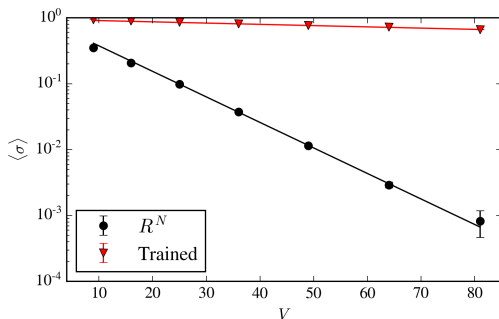
$$S = \sum_{i=1}^N \frac{m}{2} \phi_i^2 + \frac{(\phi_i - \phi_{i-1})^2}{2} + \lambda \phi_i^4, \text{ with } \lambda \in \mathbb{C}$$



- $N = 10, m = 0.5$
- 1 internal layer
- Adiabatic training

Contour deformation for complex coupling model in 1 + 1-d

The average sign over the lattice size V



- $m = 0.5, \lambda = i$
- linear contour $\tilde{\phi}(\phi) = \phi + M_R\phi + iM_I\phi$
- square lattice

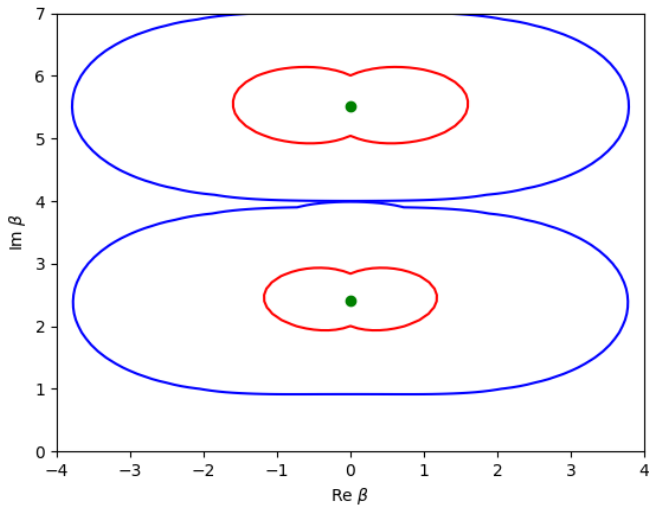
Question of existence

Do complex normalizing flows exist for any lattice field theories?

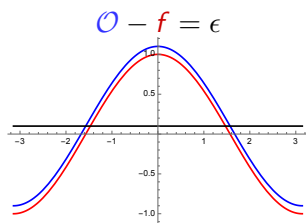
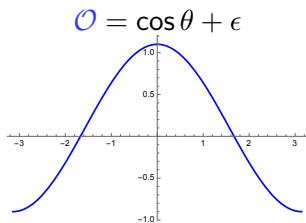
NO

(Non)-existence in $U(1)$ gauge theory in $1 + 1$ -d

$$e^{-S} = \prod_i e^{\beta \cos(\theta_i)}$$



Complex control variates



The idea is very simple...

Subtract a function f from \mathcal{O} !!

Without changing physics, so

$$\langle \mathcal{O} \rangle = \langle \mathcal{O} - f \rangle \text{ but } \text{Var}(\mathcal{O}) > \text{Var}(\mathcal{O} - f)$$

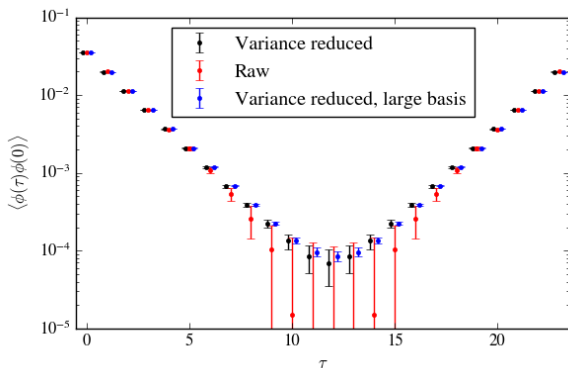
So we strictly impose

$$\langle f \rangle = \int \mathcal{D}[\phi] e^{-S(\phi)} f(\phi) = 0$$

Variance reduction for signal-to-noise problem⁵

Lattice scalar ϕ^4 theory in Euclidean

$$S = \sum_{\langle r, r' \rangle} \frac{(\phi(r) - \phi(r'))^2}{2} + \sum_r \left[\frac{m^2}{2} \phi^2(r) + \frac{\lambda}{24} \phi^4(r) \right]$$



24 × 24 lattice, $m^2 = 0.0, \lambda = 2.0$

⁵T. Bhattacharya, S. Lawrence, and J. Yoo, arXiv:2307.14950 [hep-lat]

Existence of control variates

Perfect control variates always exist!

Example:

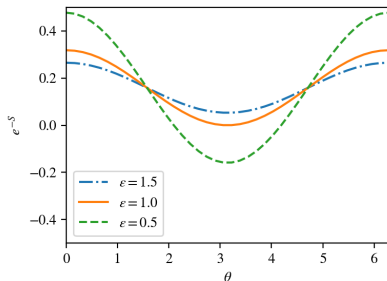
$$e^{-S(\theta; \epsilon)} = \cos(\theta) + \epsilon, \quad \theta \in [0, 2\pi)$$

What is the perfect control variates?

More generally, for any e^{-S}

$$f(x) = e^{-S(x)} - \frac{\int \mathcal{D}x e^{-S(x)}}{\int \mathcal{D}x 1}$$

(Perfect control variates are not unique)



Notes on control variates

Other strength of control variates

- Include all contour deformation methods
- No Jacobian
- Can be applied to discrete field space

How do we find good control variates?

1. Analytical (perturbative) approaches

- S. Lawrence, arXiv:2009.10901[hep-lat]
- S. Lawrence and YY, arXiv:2212.14606 [hep-lat]

2. Numerical approaches

- Start with ansatz and optimize
- Machine learning

Demonstration: Classical Ising model, Thirring model in $1 + 1$ -d

Demonstration: Classical Ising model (Lee-Yang zeros)

Classical Ising model: $S(\vec{s}) = -J \sum_{\langle i,j \rangle} s_i s_j - h \sum_i s_i$

Lee-Yang Theorem: the partition function is **0** only on with imaginary h .

Goal: Compute $Z = \sum_s e^{-S}$ at **purely imaginary magnetic field**

Measure

$$\frac{Z(h)}{Z(h=0)} = \frac{\sum_s \exp\left(J \sum_{\langle i,j \rangle} s_i s_j\right) \exp\left(h \sum_i s_i\right)}{\sum_s \exp\left(J \sum_{\langle i,j \rangle} s_i s_j\right)} = \langle e^{h \sum_i s_i} \rangle_Q$$

By replacing

$$e^{h \sum_i s_i} \rightarrow e^{h \sum_i s_i} - \mathbf{CV}$$

and optimize **CV** to minimize

$$\text{Var}\left(e^{h \sum_i s_i} - \mathbf{CV}\right)$$

Extreme learning machine

1. Prepare basis functions

$$\left\{ \sum s, \cos(\sum s), \sin(\sum s) \right\} \times s \times S(h=0)^n$$

$$(0 \leq n \leq 3)$$

2. Input basis functions to ELM

3. Take "divergence"

$$F_i = f(s_i) - f(-s_i)$$

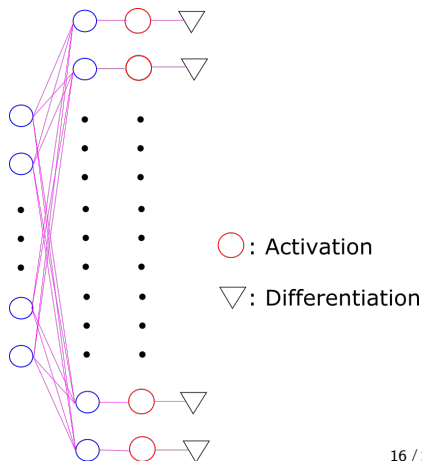
4. The $CV = \sum_i c_i F_i$

The coefficients c_i are optimized by estimating

$$M_{ij} = \langle F_i F_j \rangle, v_j = \langle \mathcal{O} F_i \rangle$$

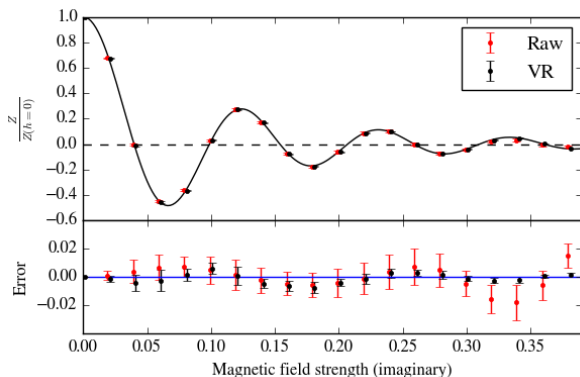
and

$$c = M^{-1} v$$



Classical Ising model⁶

At purely imaginary h , $J = 0.4 < J_c \approx 0.441$, 8×8 lattice:



- Raw: 5k samples for Z
- VR: 5k samples to optimize, 5k samples for Z

⁶S. Lawrence and YY, in preparation

Measurement of observables

Idea 1. No subtraction in the numerator

$$\langle \mathcal{O} \rangle = \frac{\int \mathcal{D}x e^{-S(x)} \mathcal{O}}{\int \mathcal{D}x e^{-S(x)}} = \frac{\int \mathcal{D}x (e^{-S(x)} - f(x)) \frac{e^{-S(x)}}{e^{-S(x)} - f(x)} \mathcal{O}}{\int \mathcal{D}x e^{-S(x)} - f(x)}$$

(Phase fluctuation moved from denominator to numerator.)

Idea 2. Subtract $\nabla \cdot (\mathcal{O} \vec{v})$ anyway

$$\begin{aligned} \langle \mathcal{O} \rangle &= \frac{\int \mathcal{D}x e^{-S(x)} \mathcal{O} - \nabla \cdot (\mathcal{O} \vec{v})}{\int \mathcal{D}x e^{-S(x)} - \nabla \cdot \vec{v}} \\ &= \frac{\int \mathcal{D}x (e^{-S(x)} - \nabla \cdot \vec{v}) \left(\mathcal{O} + \frac{\vec{v} \cdot \nabla \mathcal{O}}{e^{-S(x)} - \nabla \cdot \vec{v}} \right)}{\int \mathcal{D}x e^{-S(x)} - \nabla \cdot \vec{v}} \end{aligned}$$

Hoping that the “extra term” won't cause signal-noise problem.

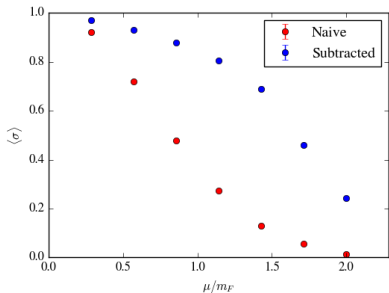
This seems to work for the density operator.... (why?)

Thirring model in 1 + 1-dimension

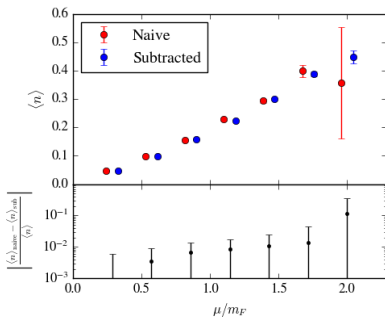
$$S = \sum_{x,\nu} \frac{2}{g^2} (1 - \cos A_\nu(x)) - \log \det K, A_\nu \in [0, 2\pi)$$

with the Dirac matrix $(\eta_0 = (-1)^{\delta_{0,x_0}}$ and $\eta_1 = (-1)^{x_0}$)

$$K[A]_{xy} = m\delta_{xy} + \frac{1}{2} \sum_{\nu=0,1} \eta_\nu e^{iA_\nu(x) + \mu\delta_{\nu,0}} \delta_{x+\nu,y} - \eta_\nu e^{-iA_\nu(y) - \mu\delta_{\nu,0}} \delta_{y+\nu,x}$$



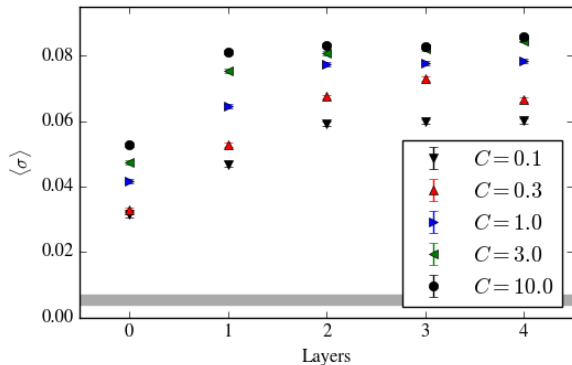
Average sign



Density

- 4×4 lattice, $m = 0.05, g = 1.0 \rightarrow m_B = 0.33(1), m_F = 0.35(2)$
- MLP with 2 inner layers

Larger networks give better vector fields



- 6×6 lattice
- $m = 0.05, g = 1.0, \mu = 0.5$

Future

- Scalable control variates for fermion sign problems
- Application of control variates to S2N problems
- Continue to check the applicability of contour deformation methods

Thank you!

Constraints on manifolds⁷

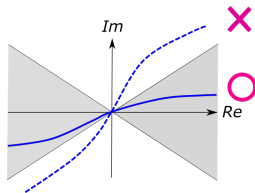
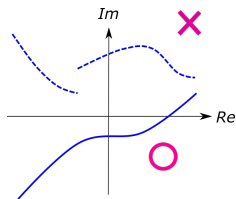
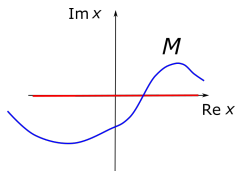
Manifolds give the correct $\langle \mathcal{O} \rangle$

$$\langle \mathcal{O} \rangle = \frac{\int_{\mathbb{R}} dy e^{-S(y)} \mathcal{O}(y)}{\int_{\mathbb{R}} dy e^{-S(y)}} = \frac{\int_{\mathcal{M}} dz e^{-S(z)} \mathcal{O}(z)}{\int_{\mathcal{M}} dz e^{-S(z)}}$$

when:

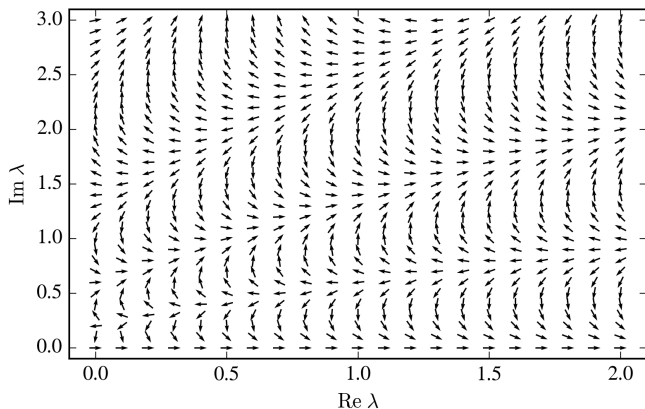
- The manifold (—) is a continuous manifold
- The manifold (—) is in “asymptotically safe” region
- Both e^{-S} and $e^{-S}\mathcal{O}$ are holomorphic functions in the region between (—) and (—)

→ **Cauchy's integral theorem!**



⁷A. Alexandru et al., Phys. Rev. D. 98, 034506(2018)

Zeros(?) of the partition function



- $m = 0.5$, 8×8 lattice
- 0 internal layer = linear transformation