Journal Club 2020/12/11 Shinichiro Akiyama

Topological phase transitions in four dimensions

N. Defenu, A. Trombettoni, and D. Zappala,

arXiv: 2003.04909 [cond-mat.quant-gas]

Abstract

4D sine-Gordon model w/ higher-derivative term may exhibit a topological phase transition analogous to the KT transition in 2D

Lifshitz point

Hornreich-Luban-Shtrikman, PRL35(1975)1678

General form of the free energy density (m: scalar order parameter)

$$f(m) = a_2 m^2 + a_4 m^4 + a_6 m^6 + \dots + c_1 (\partial m)^2 + c_2 (\partial^2 m)^2 + \dots$$

Critical point	described just by (a_2, a_4, c_1)
Tri-critical point	characterized by the a_6 term w/ the vanishing of a_4
Lifshitz point	characterized by the c_2 term w/ the vanishing of c_1

-> the Lifshitz point is analogous to a tri-critical point

Critical dimensions

- $\checkmark d=2$ is the lower critical dimension of the scalar field theory -> there is no long-range ordered phase, but the KT transition occurs
- \checkmark For the d-dim. scalar field theory, d=4 is the conjectured lower critical dimension of the Lifshitz point

Q. Is there any KT(-like) transition in 4D higher-derivative scalar field theory?

cf. Zappala, PRD98(2018)085005

Kosterlitz-Thouless transition

$$T > T_{\mathrm{KT}} \Rightarrow \mathsf{paramagnetic}$$

 $T > T_{\rm KT} \Rightarrow$ paramagnetic $T < T_{\rm KT} \Rightarrow$ vortex & anti-vortex pair

XY model (@ low T)

$$H = -J \sum_{\langle i,j \rangle} \cos(\theta_i - \theta_j) \approx \frac{J}{2} \int d\mathbf{r} (\partial \theta)^2$$

$$\checkmark \langle \cos(\theta_r - \theta_0) \rangle = \left(\frac{r}{a}\right)^{-T/2\pi J}$$
 @ $T < T_{\text{KT}}$

-> Line of fixed points

$$\to \eta = 1/4 @ T = T_{KT} (= \frac{\pi J}{2})$$

KT theory in 2D

XY model

$$H = -J \sum_{\langle i,j \rangle} \cos(\theta_i - \theta_j) \approx \frac{J}{2} \int d\mathbf{r} (\partial \theta)^2$$



vortex ⇔ Coulomb charge

Coulomb gas model

$$H = -\pi J \sum_{i \neq j} n_i n_j \log \left| \frac{r_i - r_j}{a} \right| \qquad \Longrightarrow \qquad S[\phi] = \int [(\partial \phi)^2 + g_0 (1 - \cos \beta \phi)]$$



sine-Gordon model

$$S[\phi] = \int [(\partial \phi)^2 + g_0(1 - \cos \beta \phi)]$$

KT transition @ $\beta_c^2 = 8\pi$

Coleman, PRD11(1975)2088

Outline of this study

Quartic XY model

$$H[\theta] = \frac{\mathcal{K}}{2} \int d^4 x \, \Delta \theta(x) \Delta \theta(x)$$



(ii) A possible relationship?

Coulomb gas model

$$H = -\pi J \sum_{i \neq j} n_i n_j \log \left| \frac{r_i - r_j}{a} \right|$$

Higher-derivative sine-Gordon model

$$S[\phi] = \int [(\Delta \phi)^2 + g_0(1 - \cos \beta \phi)]$$

-> (i) investigation w/ the FRG

(i) 4D higher-derivative sine-Gordon model w/ the FRG

Running action

$$\Gamma_k[\varphi] = \int d^4x \left[\frac{w_k}{2} (\Delta \varphi)^2 + V_k(\varphi) \right]$$

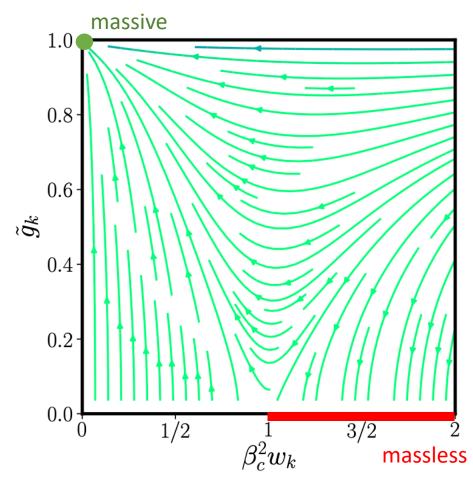
with
$$V_k(\varphi) = g_k(1 - \cos \varphi)$$

FRG flow equations

$$\partial_t w_k = -\frac{9}{160\pi^2} \frac{\tilde{g}_k^2}{(1 - \tilde{g}_k^2)^{\frac{3}{2}}}$$

$$(4 + \partial_t)\tilde{g}_k = \frac{1}{8\pi^2 w_k \tilde{g}_k} \left(1 - \sqrt{1 - \tilde{g}_k^2} \right)$$

with $t = -\log(k/\Lambda)$ and $\tilde{g}_k = g_k/k^4$



a line of attractive Gaussian fixed points for $\beta_c^2 w_k > 1$ w/ $\tilde{g}_k = 0$

$$\beta_c^2 = 64\pi^2$$
 (cf. $\beta_c^2 = 8\pi$ in 2D sine-Gordon model)

(ii) A topological configuration in 4D quartic XY model

$$H[\theta] = \frac{\mathcal{K}}{2} \int d^4x \, \Delta\theta(x) \Delta\theta(x)$$

Claim: $A_{\chi'}(x) = \frac{1}{2} \left(\alpha_4 - \frac{\pi}{2} \right) \cot \alpha_4$ gives a specific field configuration

$$\checkmark H[A_{\chi'}] = \mathcal{K}\pi^2 \ln \frac{R}{r_0}$$
 (R : large dist. cutoff, r_0 : short dist. cutoff)

$$\checkmark S[A_{\chi'}] = \ln \left[C\left(\frac{R^4}{r_0^4}\right) \right]$$

$$\checkmark \delta F = H - \mathcal{K}S = (\mathcal{K}\pi^2 - 4) \ln \frac{R}{r_0} + \text{const.}$$

$$\checkmark \mathcal{K}_c = \frac{4}{\pi^2}$$

$$x_1 - x_1' = r \sin \alpha_4 \sin \alpha_3 \sin \alpha_2$$

$$x_2 - x_2' = r \sin \alpha_4 \sin \alpha_3 \cos \alpha_2$$

$$x_3 - x_3' = r \sin \alpha_4 \cos \alpha_3$$

$$x_4 - x_4' = r \cos \alpha_4$$

(ii) Relation to the sine-Gordon model

$$H[\mathcal{G}^C] = \frac{\mathcal{K}}{2} \int d^4x \, \Delta \mathcal{G}^C(x) \, \Delta \mathcal{G}^C(x)$$

$$\mathcal{G}^{\mathcal{C}}(x) = \Sigma_i n_i \; G(x - x_i) \; \text{w/} \; n_i \in \mathbb{Z}$$

$$\Delta_x G(x - x') = -\frac{1}{(x - x')^2}$$

can be identified as the higher-derivative sine-Gordon model via ${\it Z}$

$$Z \sim \int [\mathrm{d}\phi] \exp\left[-\frac{w}{2} \int \mathrm{d}^4 r \,\Delta\phi(r) \Delta\phi(r) + 2y \int \mathrm{d}^4 r \cos\phi(r)\right]$$

$$\checkmark \frac{1}{w} \coloneqq (2\pi)^4 \mathcal{K}$$

$$\checkmark (2\pi)^4 \mathcal{K}_c = (2\pi)^4 \frac{4}{\pi^2} = 64\pi^2 (= \beta_c^2)$$

$$\checkmark \eta = \frac{1}{8\pi^2\mathcal{K}} \Rightarrow \eta = \frac{1}{32} @ \mathcal{K} = \mathcal{K}_c$$
 holds on the line of fixed points

Summary

4D systems might exhibit a topological phase analogous to the KT phase in 2D

	4D (conjecture)	2D
eta_c^2 (sine-Gordon)	$64\pi^2$	8π
\mathcal{K}_c (XY)	$4/\pi^2$	$2/\pi$
η	1/32	1/4

A delicate point is the following relationship

<u>4D quartic XY model</u>

$$H[\theta] = \frac{\mathcal{K}}{2} \int d^4 x \, \Delta \theta \Delta \theta$$

4D higher-derivative sine-Gordon model

$$H[\theta] = \frac{\mathcal{K}}{2} \int d^4 x \, \Delta \theta \Delta \theta \qquad \stackrel{?}{\Longleftrightarrow} \qquad S[\phi] = \int d^4 x [(\Delta \phi)^2 + g_0 (1 - \cos \beta \phi)]$$



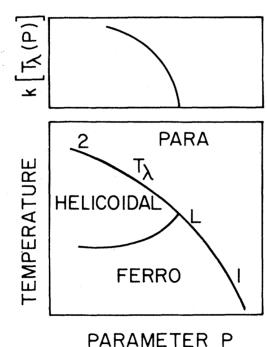
Lifshitz point

Hornreich-Luban-Shtrikman, PRL35(1975)1678

Coexistence of three phases

- ✓ homogeneous phase (order parameter is spatially uniform)
- ✓ disordered phase (order parameter is zero)
- ✓ inhomogeneous phase

 (order parameter is spatially modulated w/ a finite wave vector)



FRG flow equations

Running action

$$\Gamma_k[\varphi] = \int d^4x \left[\frac{w_k}{2} (\Delta \varphi)^2 + g_k (1 - \cos \varphi) \right]$$

$$=: V_k(\varphi)$$

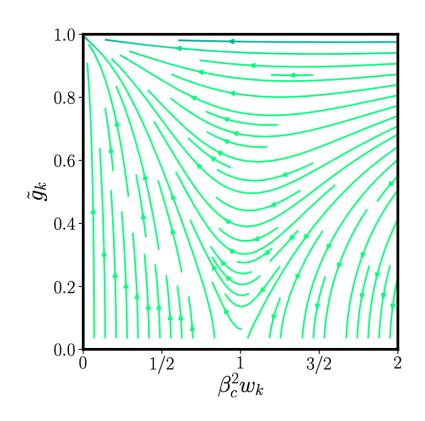
FRG flow

$$\partial_t V_k(\varphi) = \int \frac{\mathrm{d}^d q}{(2\pi)^d} G(q) \partial_t R_k(q)$$

$$\partial_t w_k = \lim_{p \to 0} \int_{-\pi}^{\pi} \frac{\mathrm{d} \varphi}{2\pi} \int \frac{\mathrm{d}^d q}{(2\pi)^d} \partial_t R_k(q)$$

$$\times G(q)^2 V_k'''(\varphi)^2 \frac{\mathrm{d}^4}{\mathrm{d} p^4} G(p+q)$$

$$G(q) = \frac{1}{w_k q^4 + V_k''(\varphi) + R_k(q)}$$



with the RG scale $t=-\log\frac{k}{\Lambda'}$, the regulator $R_k(q)=k^4$ and taking $d\to 4^+$