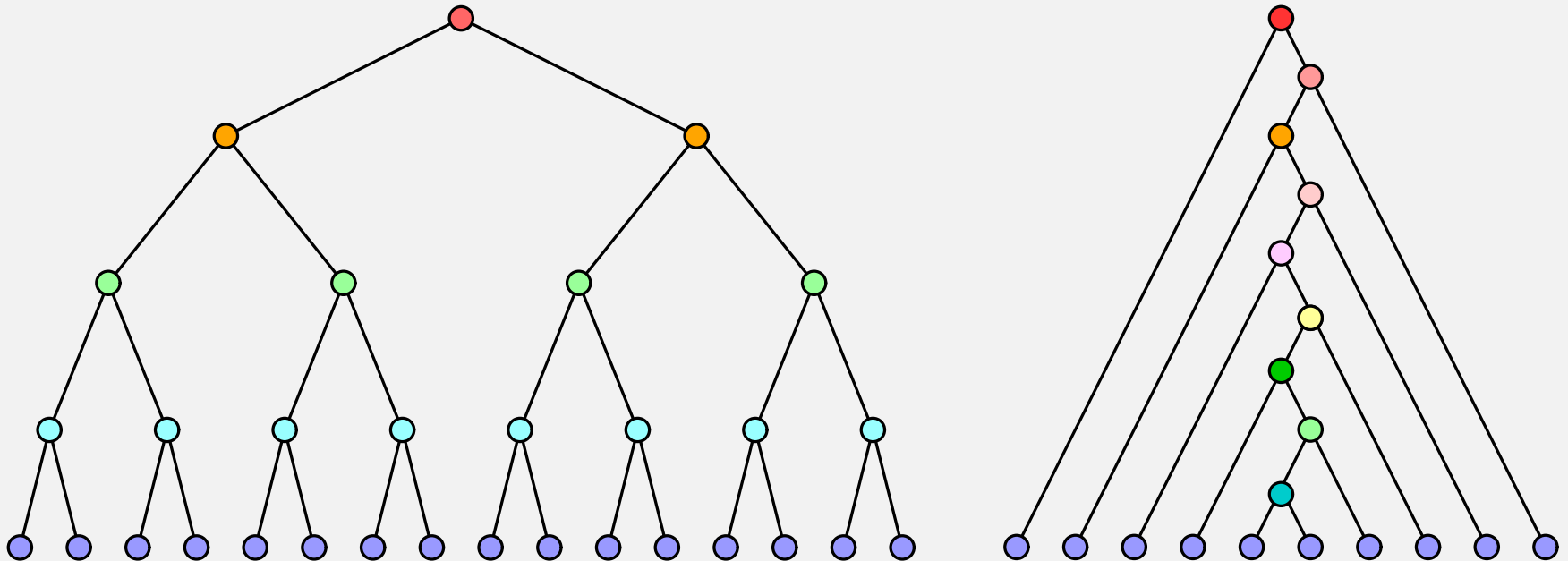


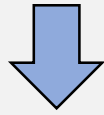
Tensor Renormalization Group Centered About a Core Tensor

Wangwei Lan and Glen Evenbly

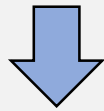
arXiv: 1906.09283 [quant-ph]



Quantum Field Theory

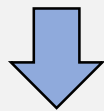


Renormalization Group
(Gell-Mann-Low Eq.)



Wilsonian RG

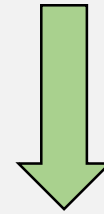
scale-dependent “effective theory”



Functional RG

**Application of
Tensor Network to QFT is
a recent hot topic!**

Critical Phenomena



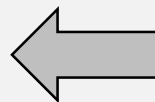
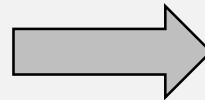
Block Spin Transformation
(Real space RG)



Density Matrix RG
1D quantum systems



**Tensor Network Scheme
Attacking higher dimension!**



Brief review of TN rep. (1/2)

Ex) 2D classical spin model (with periodic boundary)

$$Z = \sum_{\{\sigma\}} \prod_{\langle ij \rangle} \exp[-\beta K(\sigma_i, \sigma_j)]$$

Decomposition of the transfer matrix element :

$$\exp[-\beta K(\sigma_i, \sigma_j)] = \sum_l W(\sigma_i, l)W(\sigma_j, l)$$

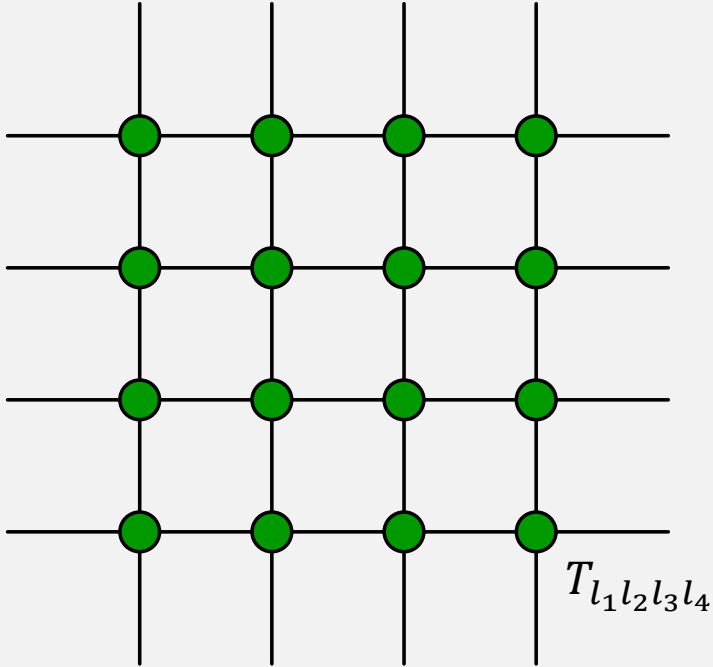
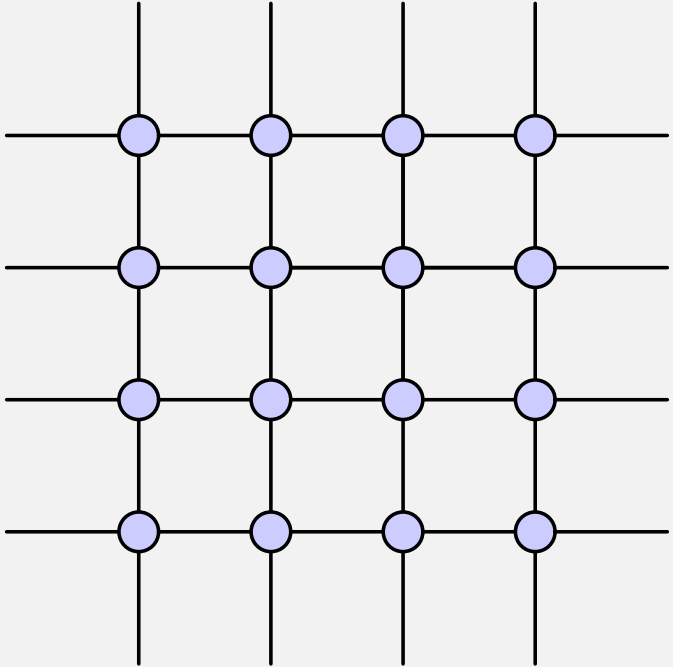
Integrating out σ 's, one obtains a local tensor;

$$T_{l_1 l_2 l_3 l_4} := \sum_{\sigma_i = \pm 1} W(\sigma_i, l_1)W(\sigma_i, l_2)W(\sigma_i, l_3)W(\sigma_i, l_4)$$

Brief review of TN rep. (2/2)

Real Space

Tensor Network
(partition function)

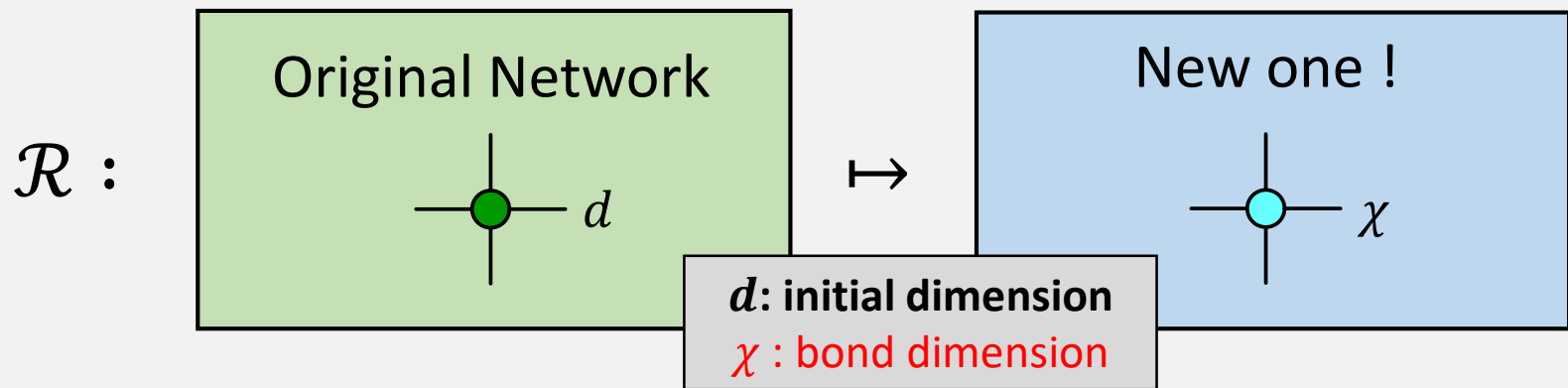
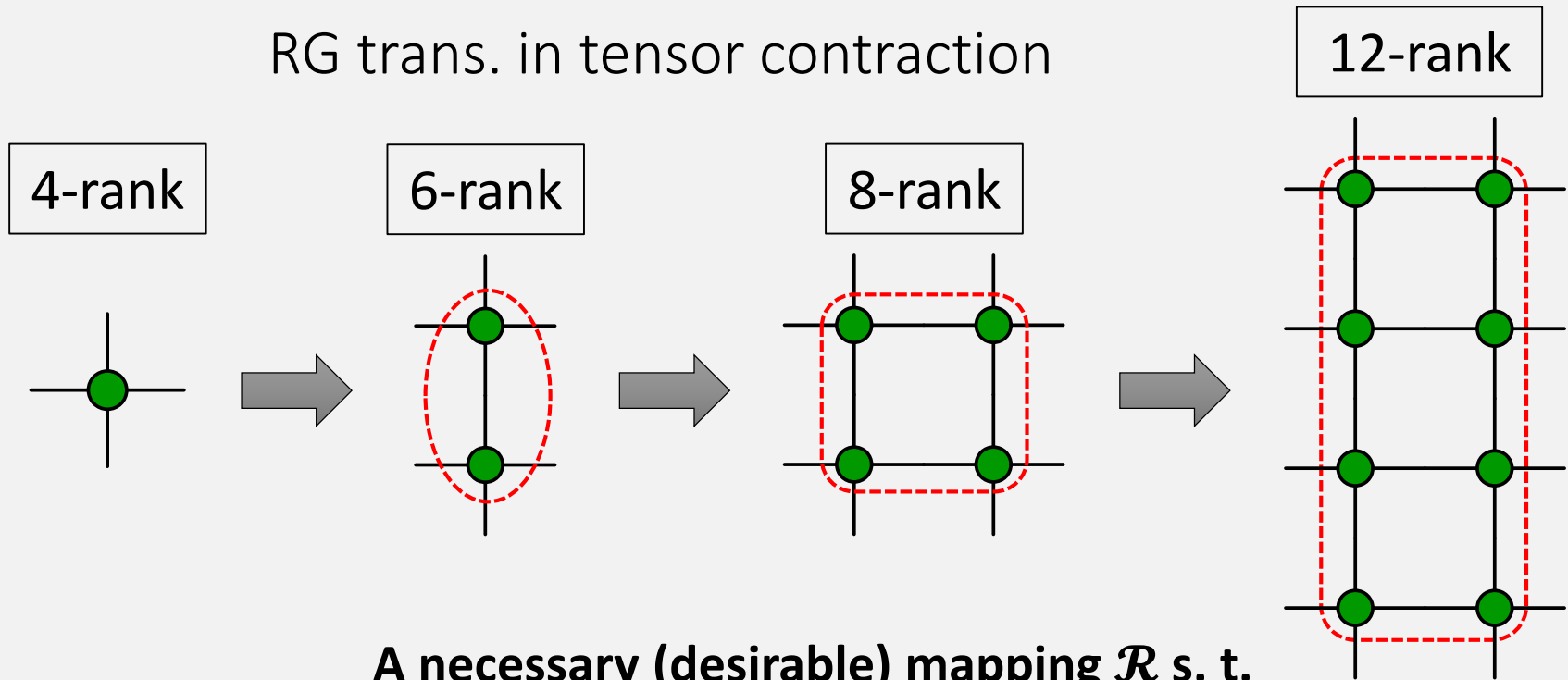


○ : spin variable

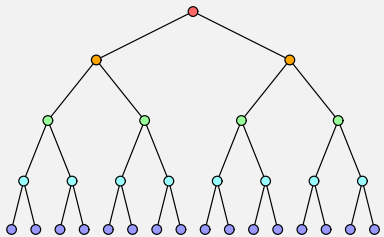
⊙ : 4-rank tensor

$$Z = \sum_{\{l\}} \prod_i T_{l_a l_b l_c l_d} =: \text{Tr} \left[\prod_i T_{l_a l_b l_c l_d} \right]$$

RG trans. in tensor contraction



Matrix decomposition technique (linear algebra) is very useful

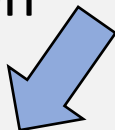


Many types of \mathcal{R} enormalization group

Tensor Renormalization Group (TRG)
 $O(\chi^6 \ln L)$ (Levin-Nave, 2007)

χ : bond dimension
 L : lattice size

Higher dimension

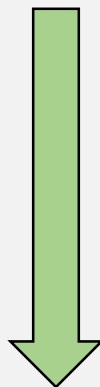


(D -dimensional) Higher-order TRG
 $O(\chi^{4D-1} \ln L)$ (Xie et al, 2012)

Accuracy



Tensor Network Renormalization
 $O(\chi^7 \ln L)$ (Evenbly-Vidal, 2015)



Cost reduction

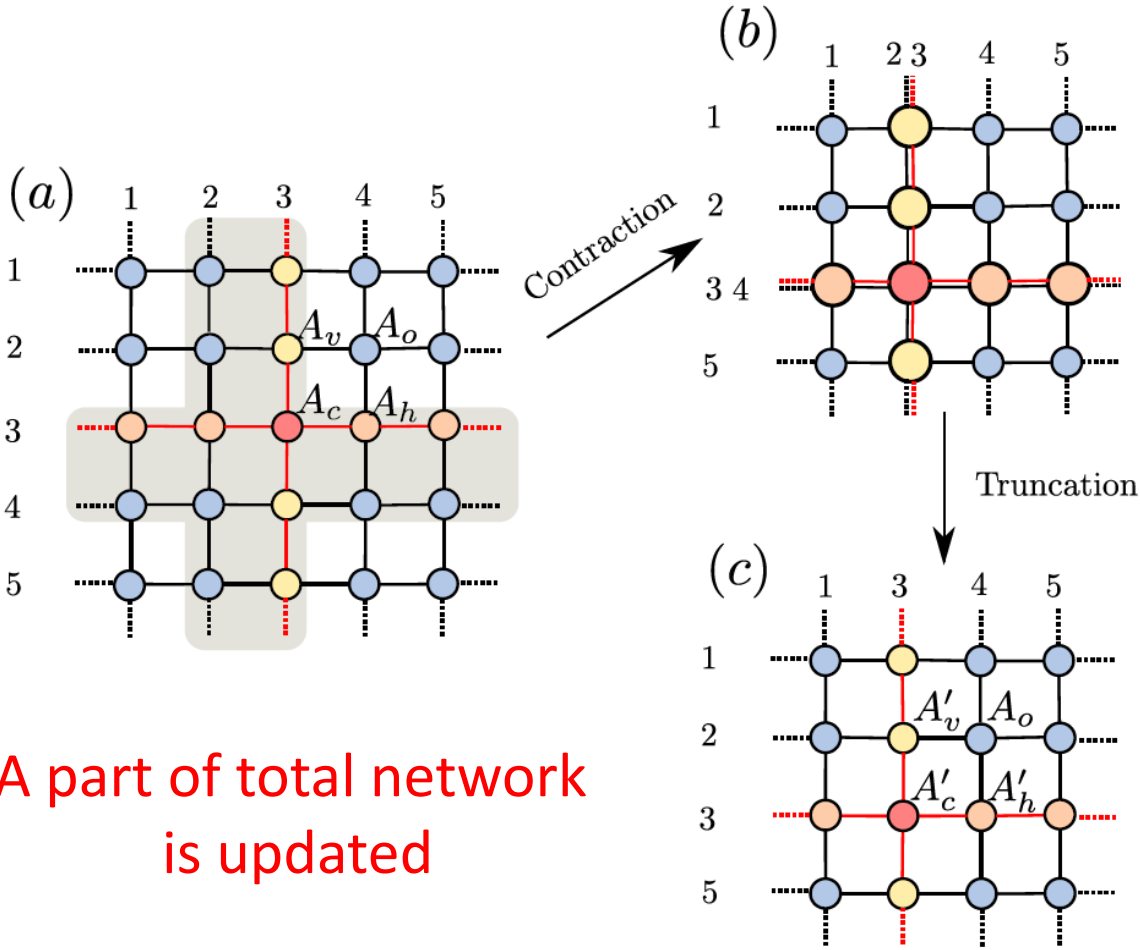
TRG with randomized SVD
 $O(\chi^5 \ln L)$ (Morita et al, 2018)

Projectively truncated TRG
 $O(\chi^5 \ln L)$ (Nakamura et al, 2019)

(D -dimensional) Anisotropic TRG
 $O(\chi^{2D+1} \ln L)$ (Adachi et al, 2019)

Core-Tensor RG
 $O(\chi^4 L)$ (Lan-Evenbly, 2019)

Algorithm of CTRG (1/3)

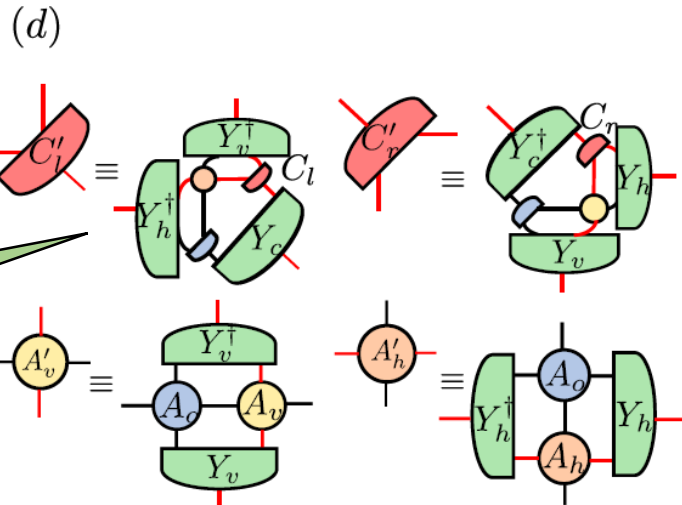
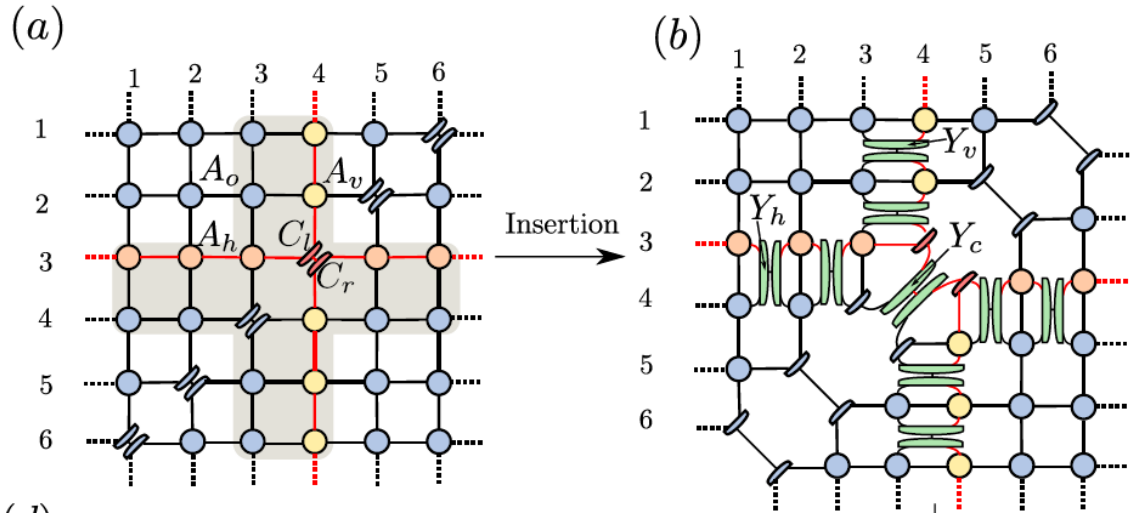


A part of total network is updated

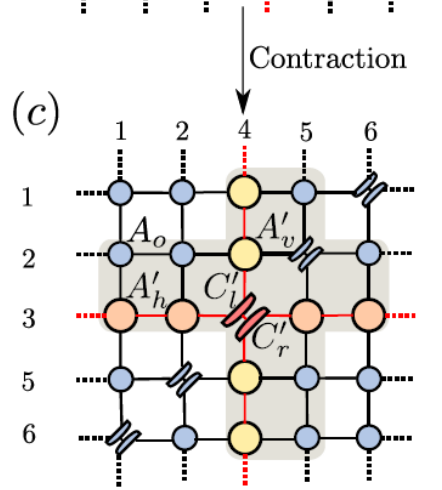
Algorithm of CTRG (2/3)

d : initial dimension
 χ : bond dimension

NOTE : Red legs run up to χ , but black ones do up to d ($\chi \gg d$)

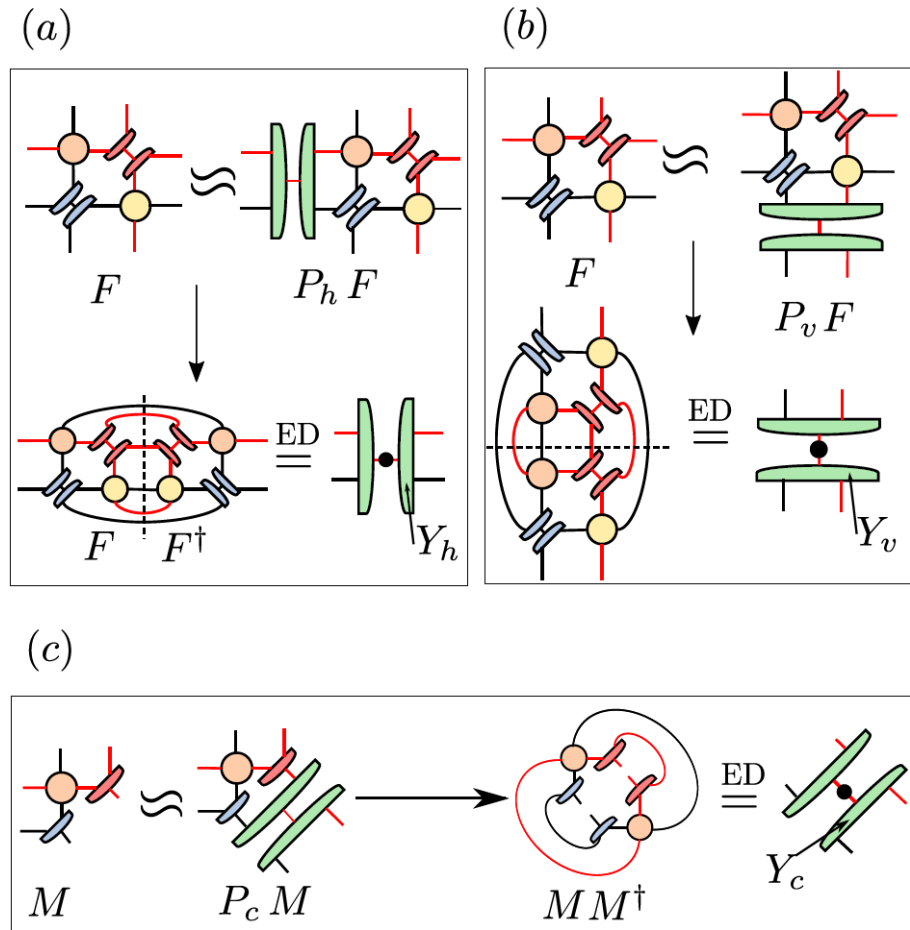


$O(d^3 \chi^4)$



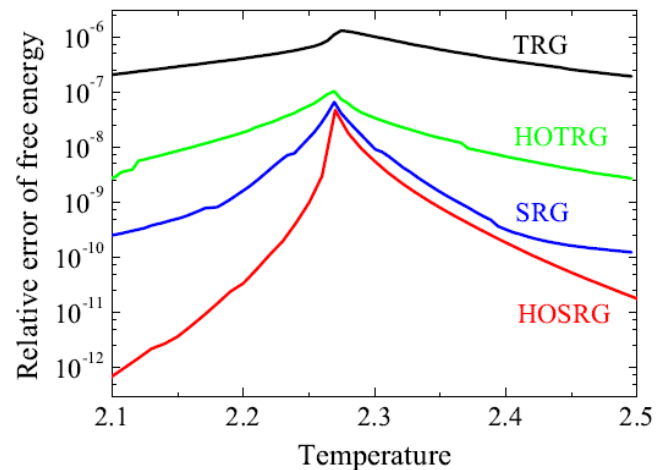
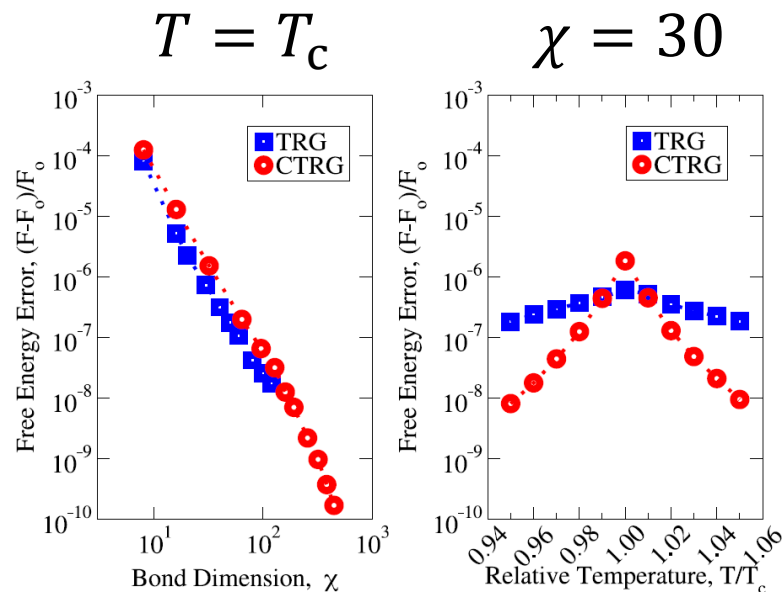
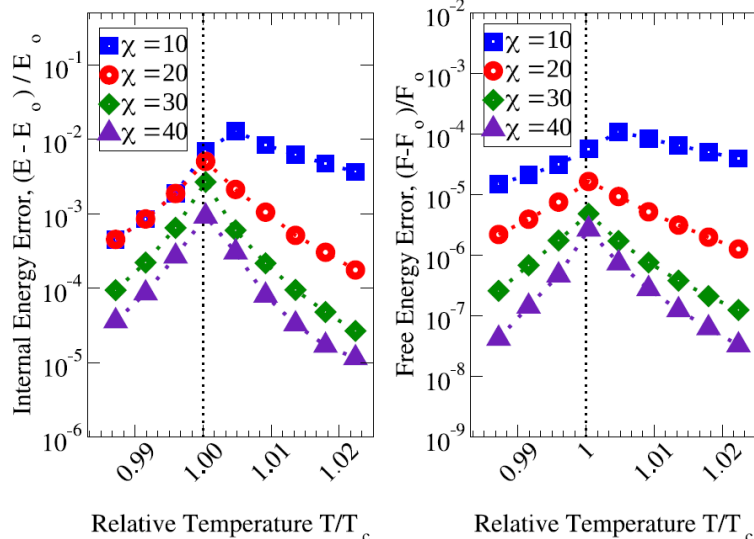
Algorithm of CTRG (3/3)

Projectors are directly decided by EVD



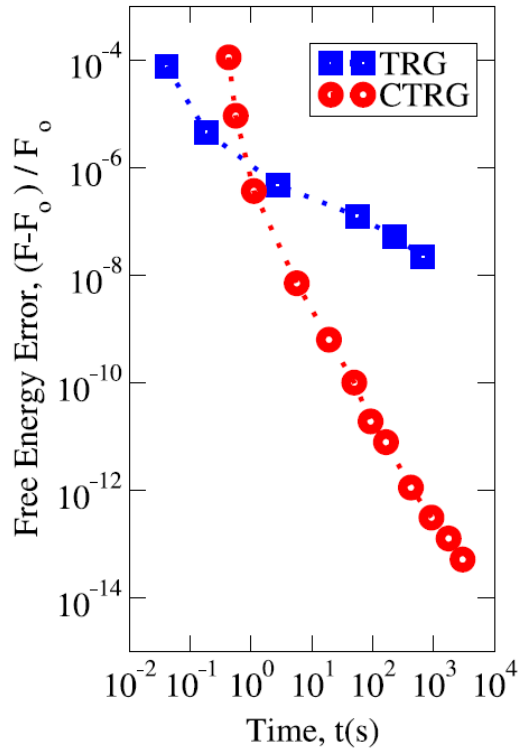
Truncation error is minimized!

Benchmark Results (2D Ising)

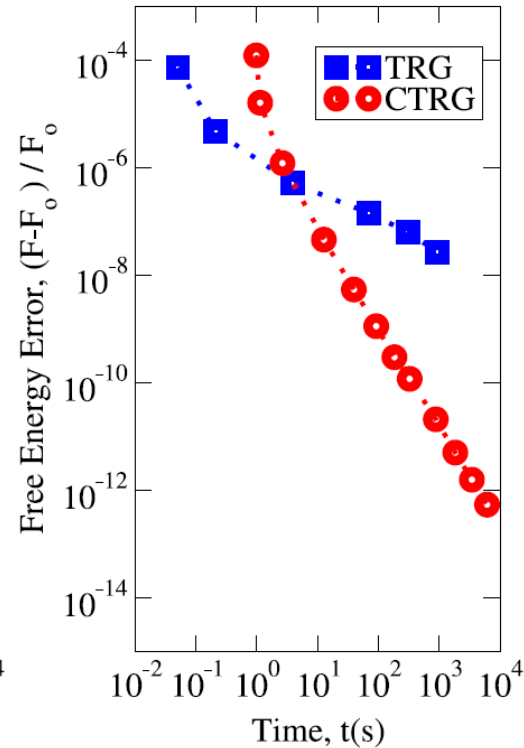


Xie et al, PRB86(2012)045139

Benchmark Results (2D Ising)



$L = 128$



$L = 256$

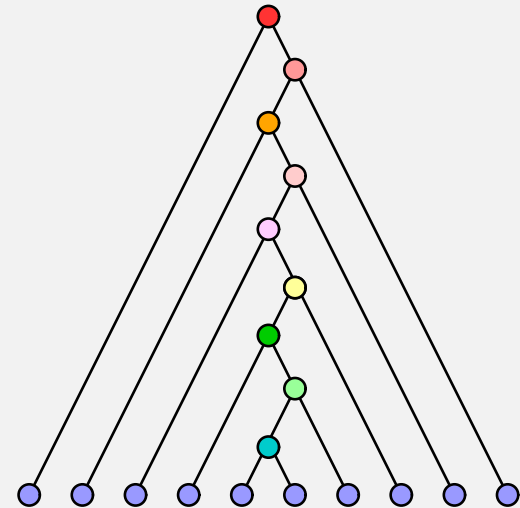
- **2 curves are crossing!** (how about comparison with HOTRG?)

Outlook

“We hypothesize that a version of CTRG generalized for higher spatial dimensions could reproduce results of equivalent accuracy to HOTRG, but with a much lower cost scaling in bond dimension χ .”

Comments

- Choice of boundary condition & Choice of $\forall L \in \mathbb{N}$
- Can be seen as a variant of CTMRG (Nishino-Okunishi, 1996), whose cost scales with $O(\chi^3 L)$



<http://quattro.phys.sci.kobe-u.ac.jp/dmrg.html>



APPENDICES

Singular Value Decomposition (SVD)

For any complex $I_1 \times I_2$ -matrix A can be written as the product

$$A = U^{(1)} S U^{(2)\dagger}$$

where

1. $U^{(1)}$ is an $I_1 \times I_1$ unitary matrix.

2. $U^{(2)}$ is an $I_2 \times I_2$ unitary matrix.

3. S is an $I_1 \times I_2$ -matrix such that

(i) Pseudo-diagonality : $S = \text{diag}(\sigma_1, \sigma_2, \dots, \sigma_{\min(I_1, I_2)})$

(ii) Ordering : $\sigma_1 \geq \sigma_2 \geq \dots \geq \sigma_{\min(I_1, I_2)} \geq 0$

σ_i 's are singular values of A and the i -th column vectors of $U^{(1)}$ and $U^{(2)}$ are, resp., i -th left and right singular vector.

Higher-Order Singular Value Decomposition (HOSVD)

Any complex $I_1 \times I_2 \times \cdots \times I_n$ -tensor A can be written as the product

$$A_{i_1 i_2 \cdots i_n} = \sum_{j_1 j_2 \cdots j_n} S_{j_1 j_2 \cdots j_n} U_{j_1 i_1}^{(1)} U_{j_2 i_2}^{(2)} \cdots U_{j_n i_n}^{(n)}$$

where

1. $U^{(k)}$ is a unitary $I_k \times I_k$ -matrix.

2. S is a complex $I_1 \times I_2 \times \cdots \times I_n$ -tensor such that

(i) Fixing the k -th index of S , say $S_{i_k=\alpha}$, and if $\alpha \neq \beta$, then

$$\sum_{i_1 i_2 \cdots i_n} S_{i_1 i_2 \cdots i_{k-1} \alpha i_{k+1} \cdots i_n} S_{i_1 i_2 \cdots i_{k-1} \beta i_{k+1} \cdots i_n} = 0$$

(ii) Ordering :

$$\|S_{i_k=\alpha}\| := \sqrt{\sum_{i_1 i_2 \cdots i_n} S_{i_1 i_2 \cdots i_{k-1} \alpha i_{k+1} \cdots i_n} S_{i_1 i_2 \cdots i_{k-1} \alpha i_{k+1} \cdots i_n}}$$

$$\|S_{i_k=1}\| \geq \|S_{i_k=2}\| \geq \cdots \geq \|S_{i_k=I_k}\| \geq 0$$

SVD introduces virtual dof

Consider the system consisting of subsystems X and Y .
Setting the pure state of the total system as

$$|\psi\rangle = \sum_{x \in X} \sum_{y \in Y} \psi(x, y) |x\rangle \otimes |y\rangle$$

If $\psi(x, y) = u(x)v(y)$, then the state is separable. Actually,

$$|\psi\rangle = \left(\sum_{x \in X} u(x) |x\rangle \right) \otimes \left(\sum_{y \in Y} v(y) |y\rangle \right)$$

SVD introduces virtual dof

Regarding $\psi(x, y)$ as a matrix element. By SVD,

$$\psi(x, y) = \sum_{l=1}^N u_l(x) \sigma_l v_l(y)$$

If $N > 1$, the state is not pure. However, as a matrix,

$$\psi = U \Sigma V^\dagger = (U \Sigma^{1/2})(V \Sigma^{1/2})^\dagger =: \tilde{U} \tilde{V}^\dagger$$

This looks very similar with $\psi(x, y) = u(x)v(y)$.