

Critical exponents of the 3d Ising and related models from Conformal Bootstrap

Ferdinando Gliozzi , Antonio Rago, arXiv:1403.6003[hep-th]

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Conformal symmetry

- At the critical point

$$\begin{aligned}x^\mu &\rightarrow e^{-\epsilon} x^\mu \\x^\mu &\rightarrow \frac{x^\mu - \epsilon^\mu x^2}{1 - 2\epsilon \cdot x + \epsilon^2 x^2}\end{aligned}$$

- Recently the critical exponents of the systems like 3D Ising model are calculated using the consequences of this symmetry.
- Gliozzi and Rago propose an alternative method to evaluate them.

Conformal symmetry

- ▶ Symmetry generators $P_\mu, M_{\mu\nu}, D, K_\mu \in SO(D, 2)$
- ▶ Primary fields

$$[M_{\mu\nu}, \mathcal{O}(0)] = \Sigma_{\mu\nu} \mathcal{O}(0)$$

$$[D, \mathcal{O}(0)] = i\Delta \mathcal{O}(0)$$

$$[K_\mu, \mathcal{O}(0)] = 0$$

classified by their scaling dimension Δ and spin l . Unitarity implies

$$\Delta \geq \begin{cases} \frac{D}{2} - 1 & (l = 0) \\ l + D - 2 & (l \geq 1) \end{cases}$$

- ▶ Other fields can be expressed as derivatives of primary fields.

Conformal symmetry

- Two and three point functions are fixed by the symmetry

$$\langle \phi(x) \phi(y) \rangle = \frac{1}{|x - y|^{2\Delta_\phi}}$$

$$\langle \phi_1(x_1) \phi_2(x_2) \phi_3(x_3) \rangle = \frac{\lambda_{123}}{|x_{12}|^{\Delta_1 + \Delta_2 - \Delta_3} |x_{23}|^{\Delta_2 + \Delta_3 - \Delta_1} |x_{13}|^{\Delta_1 + \Delta_3 - \Delta_2}}$$

$$x_{ij} = x_i - x_j$$

ϕ : primary scalar

λ_{123} : constant

OPE

- For scalar primary $\phi_{1,2}$

$$\begin{aligned}\phi_1(x) \phi_2(0) &= \sum_{\mathcal{O}} \lambda_{12\mathcal{O}} \left[\frac{1}{|x|^{\Delta_1 + \Delta_2 - \Delta_{\mathcal{O}}}} \mathcal{O}(0) + \dots \right] \\ &= \sum_{\mathcal{O}:\text{primary}} \lambda_{12\mathcal{O}} C_{12\mathcal{O}}(x, \partial) \mathcal{O}(0)\end{aligned}$$

- Four-point function

$$\begin{aligned}\langle \phi_1(x_1) \phi_2(x_2) \phi_3(x_3) \phi_4(x_4) \rangle \\ = \sum_{\mathcal{O}} \lambda_{12\mathcal{O}} \lambda_{34\mathcal{O}} C_{12\mathcal{O}}(x_1 - x_2, \partial_2) C_{34\mathcal{O}}(x_3 - x_4, \partial_4) \langle \mathcal{O}(x_2) \mathcal{O}(x_3) \rangle\end{aligned}$$

We can calculate everything once we know all the λ .

OPE

- We get the form of the four point function

$$\begin{aligned} & \langle \phi(x_1) \phi(x_2) \phi(x_3) \phi(x_4) \rangle \\ &= \sum_{\mathcal{O}} \lambda_{\phi\phi\mathcal{O}}^2 C_{\phi\phi\mathcal{O}}(x_1 - x_2, \partial_2) C_{\phi\phi\mathcal{O}}(x_3 - x_4, \partial_4) \langle \mathcal{O}(x_2) \mathcal{O}(x_3) \rangle \\ &= |x_{12}|^{-2\Delta_\phi} |x_{34}|^{-2\Delta_\phi} g(u, v) \end{aligned}$$

$$g(u, v) = 1 + \sum_{\Delta, l} \lambda_{\phi\phi\mathcal{O}}^2 G_{\Delta, l}(u, v)$$

$$u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2}, \quad v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$$

Conformal bootstrap

- There are two ways to get the four point function

$$\langle \phi(x_1) \phi(x_2) \phi(x_3) \phi(x_4) \rangle = |x_{12}|^{-2\Delta_\phi} |x_{34}|^{-2\Delta_\phi} g(u, v)$$

$$\sum_O \text{Box Diagram} = \sum_O \text{Contact Diagram}$$

- Thus we get the identity

$$v^{\Delta_\phi} g(u, v) = u^{\Delta_\phi} g(v, u)$$

$$g(u, v) = 1 + \sum_{\Delta, l} \lambda_{\phi\phi\mathcal{O}}^2 G_{\Delta, l}(u, v)$$

- We should be able to get informations on λ by studying the constraints.

Recent developments

- The identity can be rewritten as

$$\sum_{\Delta,l} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\Delta,l}(u,v) = 1$$

$$F_{\Delta,\Delta,l}(u,v) = \frac{v^{\Delta_\phi} G_{\Delta,l}(u,v) - u^{\Delta_\phi} G_{\Delta,l}(u,v)}{u^{\Delta_\phi} - v^{\Delta_\phi}}$$

- from which one gets

$$\sum_{\Delta,l} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\Delta,l}\left(\frac{1}{4}, \frac{1}{4}\right) = 1$$

$$\sum_{\Delta,l} \lambda_{\phi\phi\mathcal{O}}^2 \partial_a^{2m} \partial_b^n F_{\Delta,\Delta,l}\left(\frac{1}{4}, \frac{1}{4}\right) = 0$$

$$u = z\bar{z}, \quad v = (1-z)(1-\bar{z}), \quad z = \frac{a + \sqrt{b}}{2}$$

Recent developments

$$\sum_{\Delta,l} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\Delta,l} \left(\frac{1}{4}, \frac{1}{4} \right) = 1$$
$$\sum_{\Delta,l} \lambda_{\phi\phi\mathcal{O}}^2 \partial_a^{2m} \partial_b^n F_{\Delta,\Delta,l} \left(\frac{1}{4}, \frac{1}{4} \right) = 0$$

- ▶ Unitarity implies $\lambda_{\phi\phi\mathcal{O}}^2 > 0$.
- ▶ For 2D and 4D, one can get constraints on Δ and λ .
(Rattazzi et.al., El-Showk and Paulos)

3D Ising

- There are primary fields $\sigma, \varepsilon, \dots$ such that

$$\sigma\sigma \sim \varepsilon + \dots$$

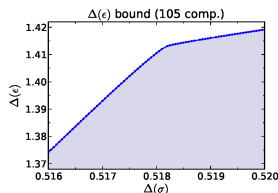
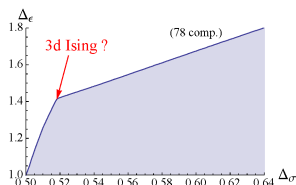
- Bootstrap equation for $\langle \sigma(x_1) \sigma(x_2) \sigma(x_3) \sigma(x_4) \rangle$:

$$\sum_{\Delta, l} \lambda_{\sigma\sigma\mathcal{O}}^2 F_{\Delta\sigma, \Delta, l} \left(\frac{1}{4}, \frac{1}{4} \right) = 1$$

$$\sum_{\Delta, l} \lambda_{\sigma\sigma\mathcal{O}}^2 \partial_a^{2m} \partial_b^n F_{\Delta\sigma, \Delta, l} \left(\frac{1}{4}, \frac{1}{4} \right) = 0$$

- For $\Delta \gg 1$, $\partial_a^{2m} \partial_b^n F_{\Delta\sigma, \Delta, l} \left(\frac{1}{4}, \frac{1}{4} \right) \propto r^\Delta P(\Delta)$ ($r < 1$) with P a polynomial.
- One can approximate these equations by finite set of (Δ, l) (~ 200) and get an upper bound for Δ_ε .

3D Ising



- ▶ From the values of $\Delta_\sigma, \Delta_\epsilon$ the 3D Ising model is at the cusp of the graph.
- ▶ Assuming it is exactly at the cusp, one can get various quantities very accurately. (El-Showk et.al.)

$$\Delta_\sigma = 0.518154(15), \Delta_\epsilon = 1.41267(13), \text{ etc.}$$

Proposal of Gliozzi and Rago

- Start from the bootstrap equations

$$\sum_{\Delta, l} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta, \phi, \Delta, l} \left(\frac{1}{4}, \frac{1}{4} \right) = 1$$

$$\sum_{\Delta, l} \lambda_{\phi\phi\mathcal{O}}^2 \partial_a^{2m} \partial_b^n F_{\Delta, \phi, \Delta, l} \left(\frac{1}{4}, \frac{1}{4} \right) = 0$$

and the fusion rule (first few terms)

$$[\Delta_\phi] \times [\Delta_\phi] = \sum_i [\Delta_i, l_i]$$

- Prepare N set of (Δ, l) and consider $M (> N)$ such equations. From the condition that the linear equations have a solution, we get Δ .

Results 1

Lee-Yang edge singularity

$$[\Delta_\phi] \times [\Delta_\phi] = 1 + [\Delta_\phi] + [D, 2] + [\Delta_4, 4] + [\Delta'] + \dots$$

<i>D</i> dimensional Yang-Lee model—the edge exponent σ					
<i>D</i>	bootstrap	Ising in <i>H</i>	Fluids	Animals	ϵ -expansion
2	-0.1664(5)	-0.1645(20)	-0.161(8)	-0.165(6)	(exact -1/6)
3	0.085(1)	0.077(2)	0.0877(25)	0.080(7)	0.079-0.091
4	0.2685(1)	0.258(5)	0.2648(15)	0.261(12)	0.262-0.266
5	0.4105(5)	0.401(9)	0.402(5)	0.40(2)	0.399-0.400
6	0.4999(1)	0.460(50)	0.465(35)	—	1/2

- One does not need unitarity to apply this method.

Results 2

3D Ising

► fusion rule

$$\begin{aligned} [\Delta_\sigma] \times [\Delta_\sigma] &= 1 + [\Delta_\varepsilon] + [\Delta_{\varepsilon'}] + [\Delta_{\varepsilon''}] + [\Delta_{\varepsilon'''}] \\ &\quad + [3, 2] + [\Delta_4, 4] + [\Delta_6, 6] + \cdots \end{aligned}$$

► $\Delta_4 = 5.0208(12)$

scalar operators	σ	σ'	ε	ε'	ε''	ε'''
Δ , best estimates	0.51813(5)	$\gtrsim 4.5$	1.41275(25)	3.84(4)	4.67(11)	—
Δ , bootstrap	0.5171(1)	4.05(5)	1.413(2)	3.79(3)	4.60(7)	5.80(4)

spinning operators	$[\Delta_2, 2]$	$[\Delta'_2, 2]$	$[\Delta_4, 4]$	$[\Delta'_4, 4]$	$[\Delta_6, 6]$
Δ , best estimates	—	—	5.0208(12)	—	7.028(8)
Δ , bootstrap	5.117(1)	6.20(1)	input	6.70(1)	7.04(2)