

A proposal for B-physics on current lattices

B.Blossier, et al.[ETMC Collaboration], JHEP 1004(2010) 049

P.Dimopoulos et al.[ETMC Collaboration], JHEP01(2012) 046

PoS LATTICE2012 (2012) 104 (1211.0568)

PoS LATTICE2012 (2012) 105 (1211.0565)

N.Carrasco et al.[ETMC Collaboration], arXiv:1308.1851

PoS LATTICE2013(2013) (1311.2837)

B-physics on the lattice

- Precise SM predictions need good control of systematic errors on the lattice.
- discretization error $am \sim 1$ in heavy quark sector $O(ma)$
 - A. lattice HQET and lattice NRQCD
 - B. Relativistic heavy quark action
(Fermilab, Tsukuba, RHQ, Oktay-Kronfeld)
 - C. light quark action

light quark action

- fine lattice $ma < 1$ - brute force
- improvement to suppress systematic errors
 $O(ma)$
 - A. Clover improvement(Wilson)
 - B. HISQ(High Improved Staggered Quark)
- “Ratio method” with the improved action
[JHEP 1004(2010) 049]

Ratio method

[JHEP 1004(2010) 049]

- interpolation available data (around charm quark region) and static limit of Heavy quark effective theory

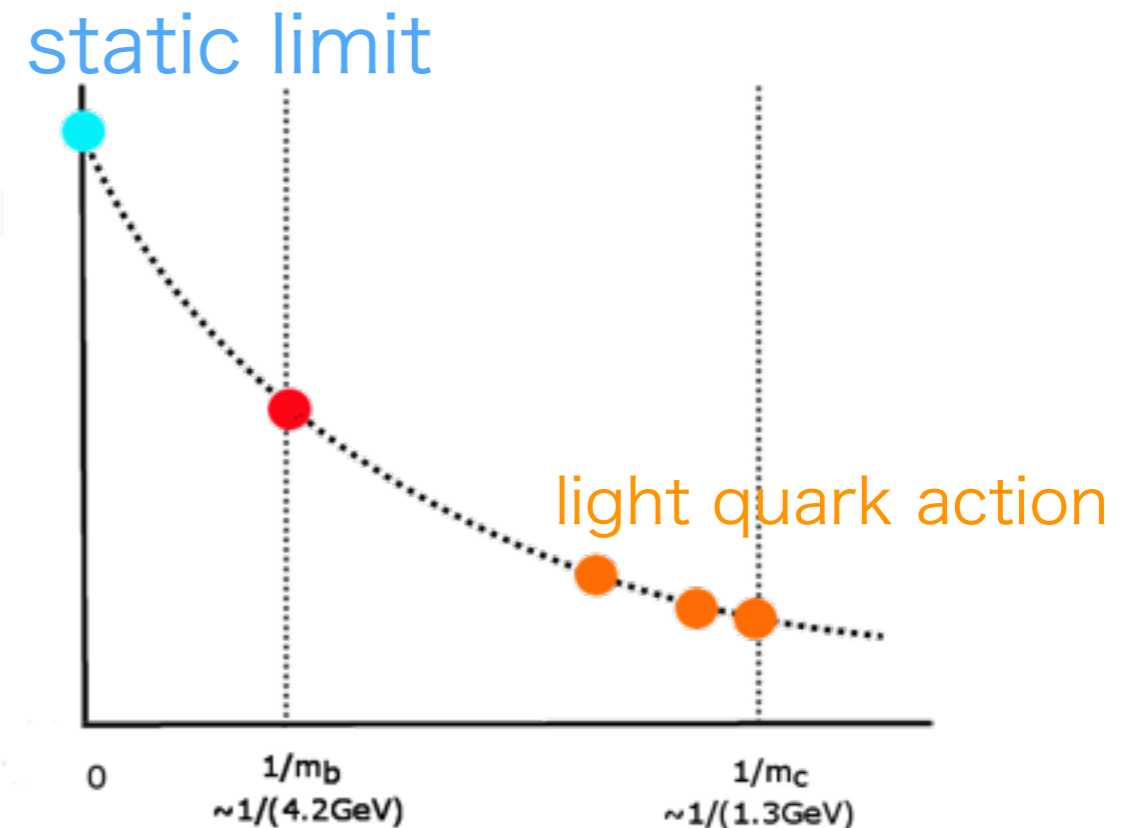
- Ratio for heavy-light observable - mild dependence of quark mass and cutoff effects.

leading order of HQET \Rightarrow static limit

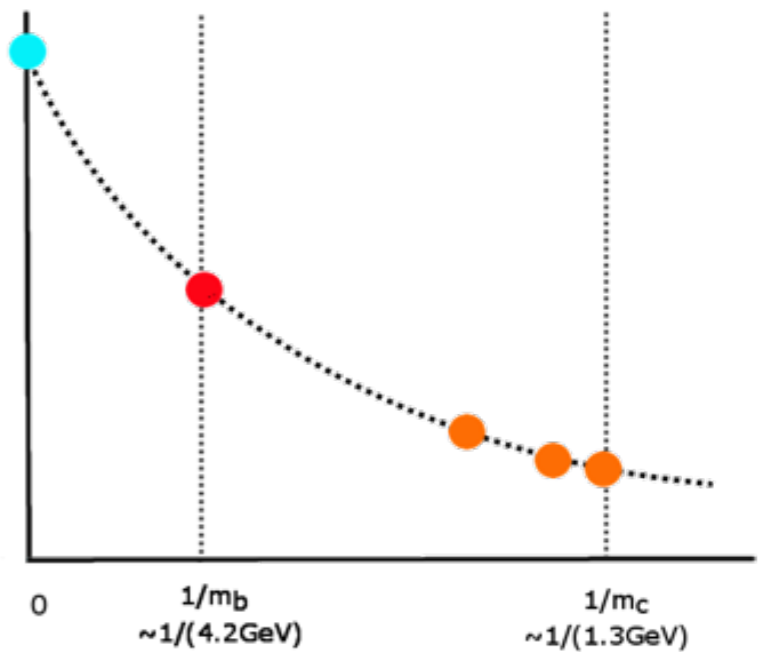
- Twisted mass Wilson fermion

$O(a)$ -improved action

- b-quark mass
- decay constant



Ratio method for the b-quark mass



- HQET prediction in static limit

$$\lim_{\mu_h^{pole} \rightarrow \infty} \left(\frac{M_{hl}}{\mu_h^{pole}} \right) = const$$

M_{hl} : heavy – light PS meson mass
 μ_h^{pole} : heavy quark pole mass

- Consider heavy quark mass set $\hat{\mu}_h^{(1)} < \hat{\mu}_h^{(2)} < \dots < \hat{\mu}_h^{(n)}$
 $\hat{\mu}_h^{(i)}$: renormalized at 2.0GeV in MS bar scheme.

- lattice ratios

$$y^L \left(\hat{\mu}^{(n)}; \hat{\mu}_l, a \right) = \frac{M_{hl}^L \left(\hat{\mu}_h^{(n)}; \hat{\mu}_l, a \right)}{M_{hl}^L \left(\hat{\mu}_h^{(n-1)}; \hat{\mu}_l, a \right)} \cdot \frac{\rho(\hat{\mu}_h^{(n-1)}) \hat{\mu}_h^{(n-1)}}{\rho(\hat{\mu}_h^{(n)}) \hat{\mu}_h^{(n)}} \xrightarrow[\substack{\hat{\mu}_l \rightarrow \hat{\mu}_{u/d} \\ a \rightarrow 0}]{n \rightarrow \infty} 1$$

$$\rho(\hat{\mu}_h) \hat{\mu}_h = \mu_h^{pole} \quad \rho(\hat{\mu}_h) : \text{transform factor}$$

This factor is known up to NNLL order in continuum perturbation theory.[hep-ph/9912391]

Ratio method

for the b-quark mass

- quark mass sequence with fixed ratio

$$\lambda = \frac{\hat{\mu}_h^{(n)}}{\hat{\mu}_h^{(n-1)}} = \frac{x^{(n-1)}}{x^{(n)}} \quad x^{(n)} = \frac{1}{\hat{\mu}_h^{(n)}}$$

- y ratio with λ

$$\begin{aligned} y \left(x^{(n)}, \lambda; \hat{\mu}_{u/d} \right) &\equiv \lim_{\hat{\mu}_l \rightarrow \hat{\mu}_{u/d}} \lim_{a \rightarrow 0} y^L \left(x^{(n)}, \lambda; \hat{\mu}_l \right) \\ &= \lambda^{-1} \frac{M_{hu/d} \left(1/x^{(n)} \right)}{M_{hu/d} \left(1/\lambda x^{(n)} \right)} \cdot \frac{\rho(\lambda x^{(n)})}{\rho(x^{(n)})} \end{aligned}$$

- ansatz for y

$$y \left(x, \lambda; \hat{\mu}_{u/d} \right) \Big|_p = 1 + \eta_1 x + \eta_2 x^2 \quad \lim_{x \rightarrow 0} y \left(x^{(n)}, \lambda; \hat{\mu}_{u/d} \right) = 1$$

ρ up to $N^{p-1}L$ order

Simulation details

- Tree-level Symanzik gauge action
- Nf=2 dynamical configurations
- Twisted mass fermion action at maximal twist
- Renormalization constant are determined by RI-MOM scheme

$$\hat{\mu}_h^{(1)} = 1.230\text{GeV}, \hat{\mu}_h^{(2)} = \lambda \hat{\mu}_h^{(1)} = 1.572\text{GeV}, \hat{\mu}_h^{(3)} = \lambda^2 \hat{\mu}_h^{(2)} = 2.009\text{GeV}, \hat{\mu}_h^{(4)} = \lambda^3 \hat{\mu}_h^{(3)} = 2.568\text{GeV} \quad \lambda = 1.278$$

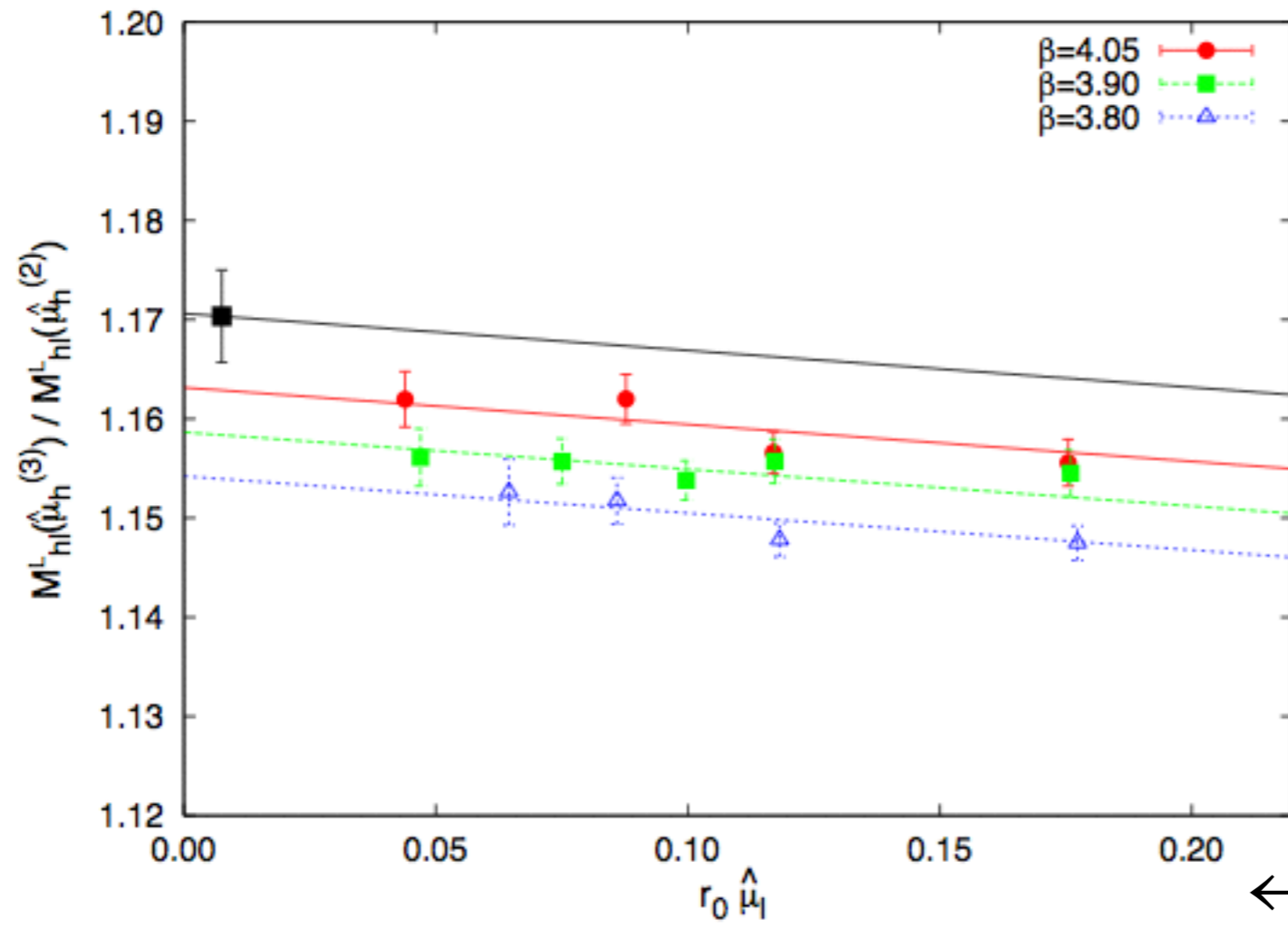
β	$a^{-4}(L^3 \times T)$	$a\mu_\ell = a\mu_{sea}$	$a\mu_s$	$a\mu_h$
3.80	$24^3 \times 48$	0.0060, 0.0080	0.0200, 0.0250	0.2700, 0.3100
		0.0110, 0.0165	0.0300, 0.0360	0.3550, 0.4350
3.90	$24^3 \times 48$	0.0040, 0.0064	0.0220, 0.0270	0.2500, 0.3200
		0.0085, 0.0100	0.0320	0.3900, 0.4600
		0.0150		
3.90	$32^3 \times 64$	0.0030, 0.0040	0.0220, 0.0270	0.2500, 0.3200
4.05	$32^3 \times 64$	0.0030, 0.0060	0.0150, 0.0180	0.2000, 0.2300
		0.0080, 0.0120	0.0220, 0.0260	0.2600, 0.3150

cutoff scale

$1/a = 1.7 - 2.6 \text{ GeV}$

measured correlator = 240

cutoff-scale and light quark mass dependence

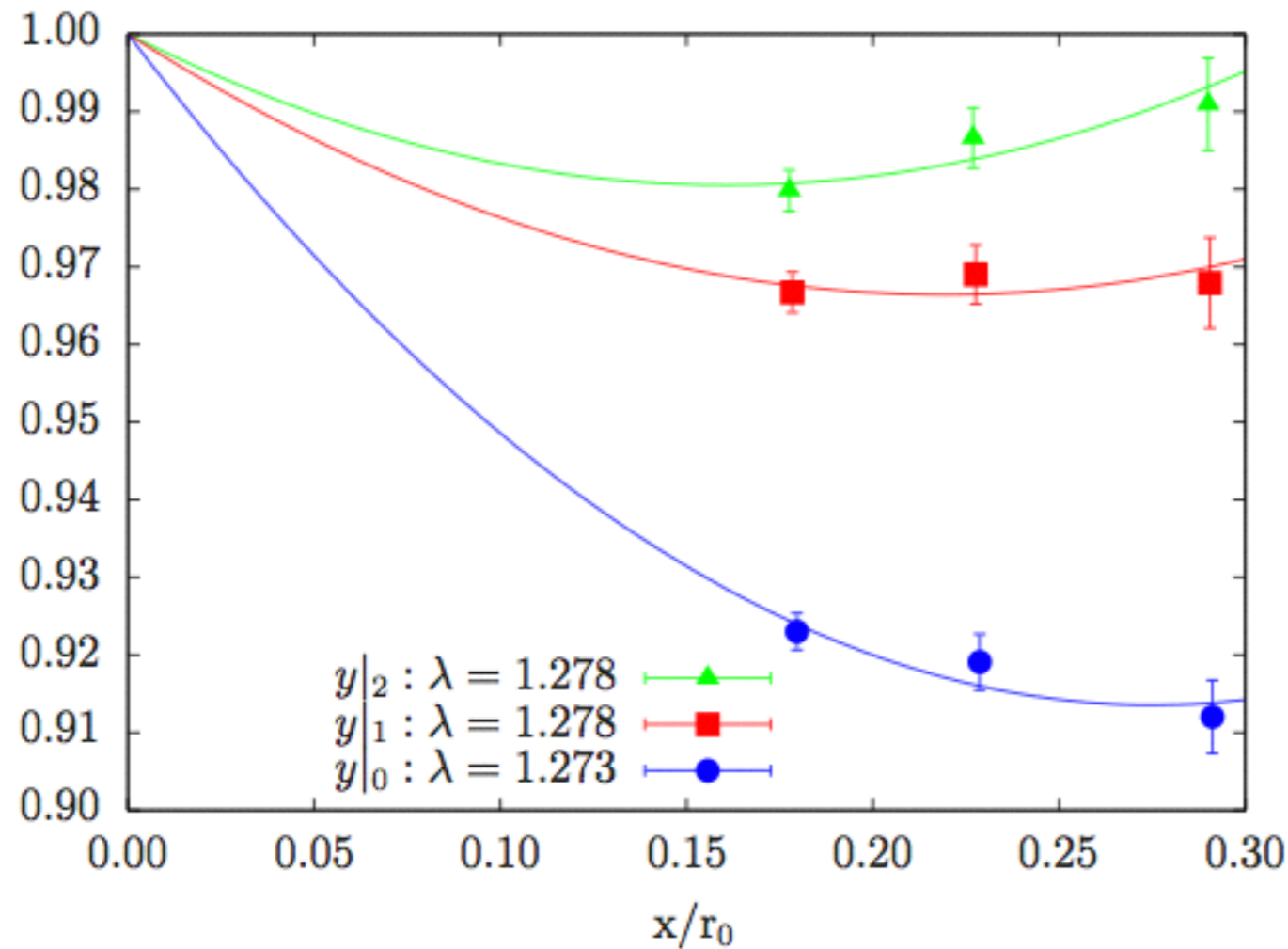


extrapolation
with HMChPT

← light quark

Figure 1: Lattice spacing and $\hat{\mu}_\ell$ dependence of the ratio $M_{hl}^L(\hat{\mu}_h^{(3)}; \hat{\mu}_\ell, a) / M_{hl}^L(\hat{\mu}_h^{(2)}; \hat{\mu}_\ell, a)$. The black square with its error is the combined continuum and chirally ($\hat{\mu}_\ell \rightarrow \hat{\mu}_{u/d}$) extrapolated value. Here and in all the following figures uncertainties possibly affecting the value of the variable in the horizontal axis are propagated to the quantity plotted on the vertical axis.

y ratios from current data



$$\frac{\rho(\hat{\mu}_h^{(n-1)})\hat{\mu}_h^{(n-1)}}{\rho(\hat{\mu}_h^{(n)})\hat{\mu}_h^{(n)}}$$

$$y(x, \lambda; \hat{\mu}_{u/d})|_p = 1 + \eta_1 x + \eta_2 x^2$$

➔ η_1, η_2

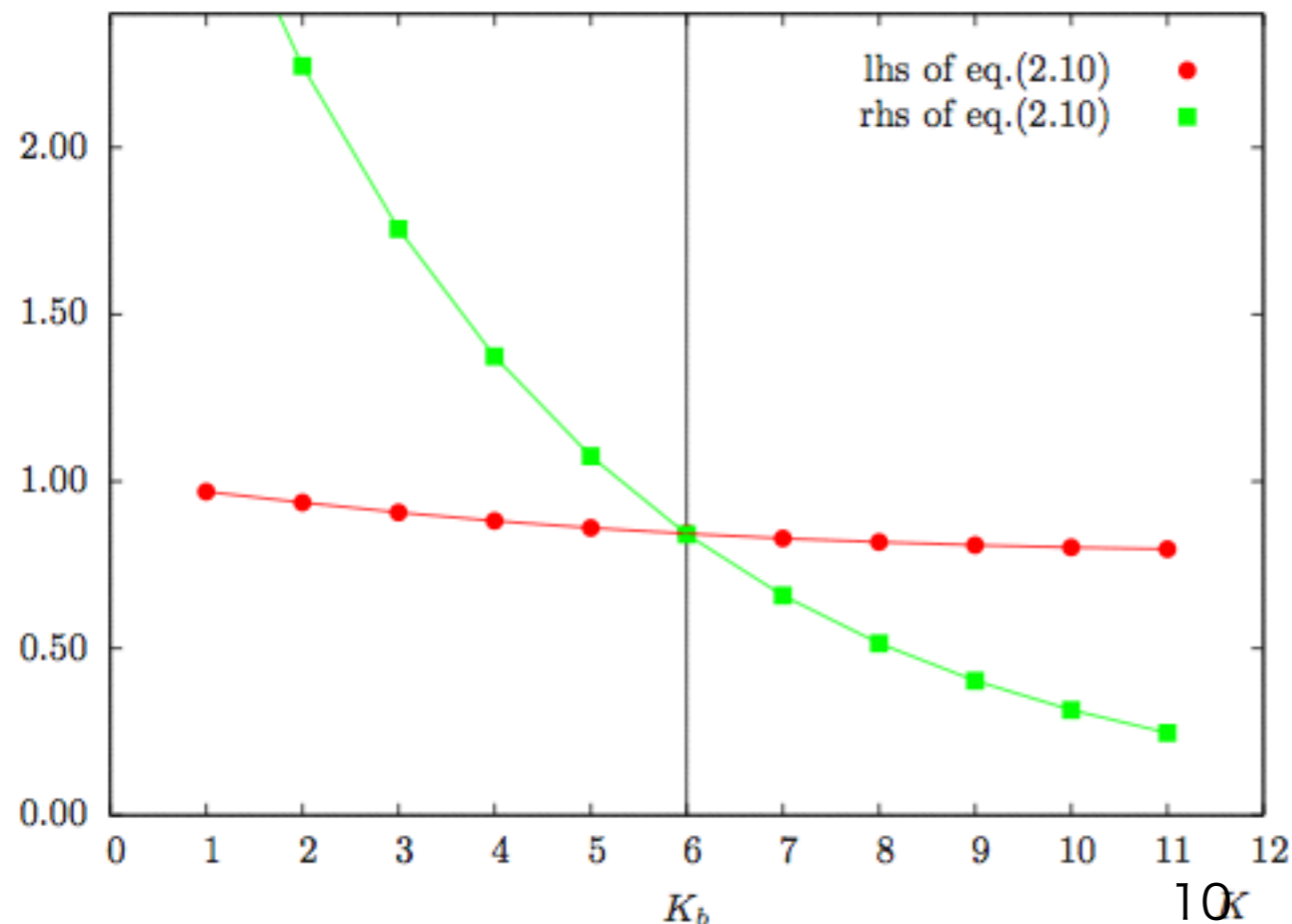
Figure 2: Continuum data for $y|_0$ (blue dots), $y|_1$ (red squares), $y|_2$ (green triangles). The corresponding best fit curves are drawn with $\lambda = 1.273$ (lower curve, in blue) and $\lambda = 1.278$ (middle curve, in red and upper curve, in green). In all cases $\mu_\ell \rightarrow \mu_{u/d}$.

determination of B-quark mass

- A following chain equation beyond current data

$$y_p^{(2)} y_p^{(3)} \dots y_p^{(K+1)} = \lambda^{-K} \frac{M_{hu/d}^{exp} \left(\hat{\mu}_h^{(K+1)} \right)}{M_{hu/d} \left(\hat{\mu}_h^{(1)} \right)} \cdot \left[\frac{\rho(\hat{\mu}_h^{(1)})}{\rho(\hat{\mu}_h^{(K+1)})} \right]_p$$

$\hat{\mu}_b = \lambda^{K_b} \hat{\mu}^{(1)}$



$$\hat{\mu}_b^{(\overline{MS}, N_f=2)} = 4.63(27) GeV$$

$$\hat{\mu}_b^{(\overline{MS}, PDG2014)} = 4.18(3) GeV$$

sea s,c quark effect

$$\hat{\mu}_b^{(\overline{MS}, N_f=4)} = 4.29(13) GeV$$

[N.Carrasco et al, arXiv:1311.2837]

Discussion and error budget

$$y_p^{(2)} y_p^{(3)} \dots y_p^{(K+1)} = \lambda^{-K} \frac{M_{hu/d}(\hat{\mu}_h^{(K+1)})}{M_{hu/d}(\hat{\mu}_h^{(1)})} \cdot \left[\frac{\rho(\hat{\mu}_h^{(1)})}{\rho(\hat{\mu}_h^{(K+1)})} \right]_p$$

- reliability of truncation of ρ (NL order)

L order $\sim 2\%$

NNL order $\sim 1\%$

- Pseudo-scalar meson mass in charm region

5% (scale setting, renormalization constant, ...)

decay constant : f_B, f_{B_s}

- Precise determination of PS decay constant f_B, f_{B_s} is significant for “New Physics” sensitive processes e.g. $B \rightarrow \tau \nu$, $B_s \rightarrow \mu^+ \mu^-$.
- ratio method for f_B, f_{B_s}

$$\lim_{\hat{\mu}_h^{pole} \rightarrow \infty} f_{hl} \sqrt{\mu_h^{pole}} = \text{constant}$$

- A very similar strategy with determination of b quark mass

ratio method

for decay constant

- lattice ratio

$$z(x, \lambda; \hat{\mu}_l) = \lambda^{1/2} \frac{f_{hl}(1/x)}{f_{hl}(1/\lambda x)} \cdot \frac{C_A^{stat}(\lambda x) [\rho(x)]^{1/2}}{C_A^{stat}(x) [\rho(\lambda x)]^{1/2}}$$

$$\Phi_{hl}(\mu^*) = [C_A^{stat}(\mu^*, \hat{\mu}_h)]^{-1} \cdot \Phi_{hl}^{QCD}(\hat{\mu}_h) \quad \text{up to 3 loops [hep-ph/0303113]}$$

- f_B : light \rightarrow u/d, f_{Bs} : light \rightarrow s

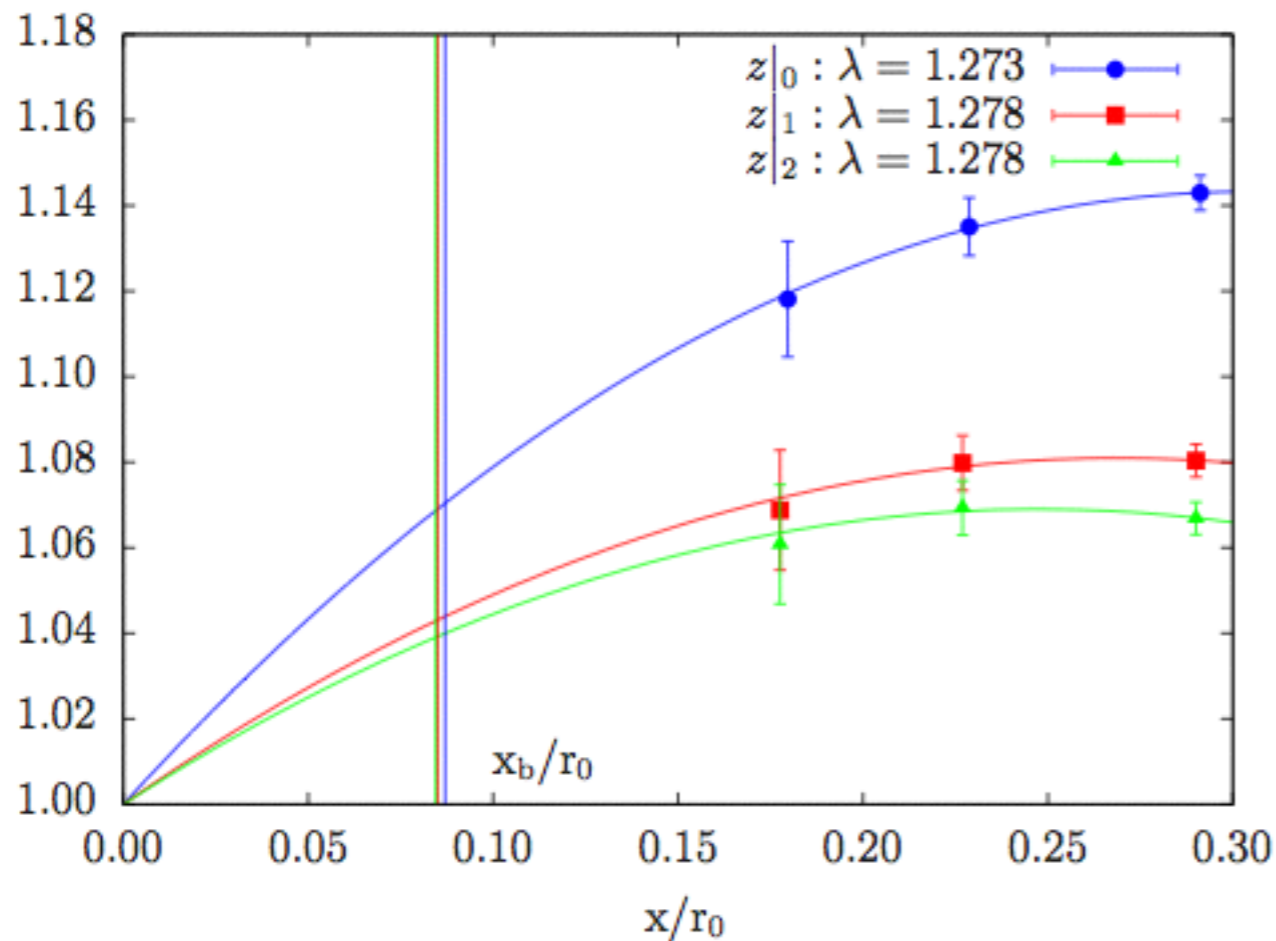
- asymptotic behavior z $\lim_{x \rightarrow 0} z(x, \lambda; \hat{\mu}_{u/d}) = 1$

- ansatz

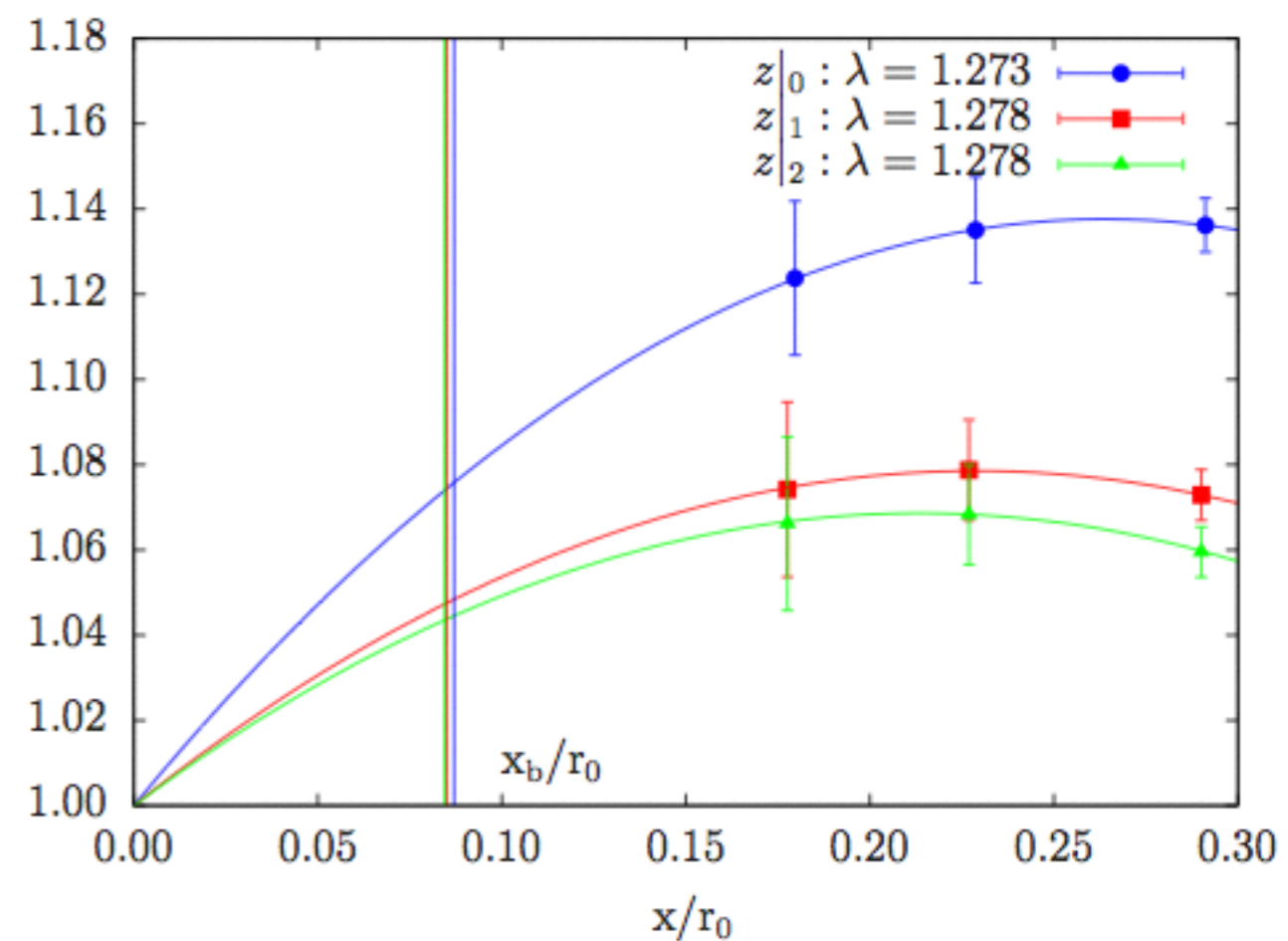
$$z(x, \lambda; \hat{\mu}_{u/d})|_p = 1 + \zeta_1 x + \zeta_2 x^2 \quad \longrightarrow \quad \zeta_1, \zeta_2$$

interpolation to b quark region

$$z(x, \lambda; \hat{\mu}_{u/d})|_p = 1 + \zeta_1 x + \zeta_2 x^2$$



light - u/d



light - s

determination of f_B and f_{B_s}

- the iterative formula

$$z_p^{(2)} z_p^{(3)} \dots z_p^{(K+1)} = \lambda^{K/2} \frac{f_{hl}(\hat{\mu}_h^{(K+1)})}{f_{hl}(\hat{\mu}_h^{(1)})} \cdot \left[\frac{C_A^{stat}(\hat{\mu}_h^{(1)})}{C_A^{stat}(\hat{\mu}_h^{(K+1)})} \sqrt{\frac{\rho(\hat{\mu}_h^{(K+1)})}{\rho(\hat{\mu}_h^{(1)})}} \right]_p$$

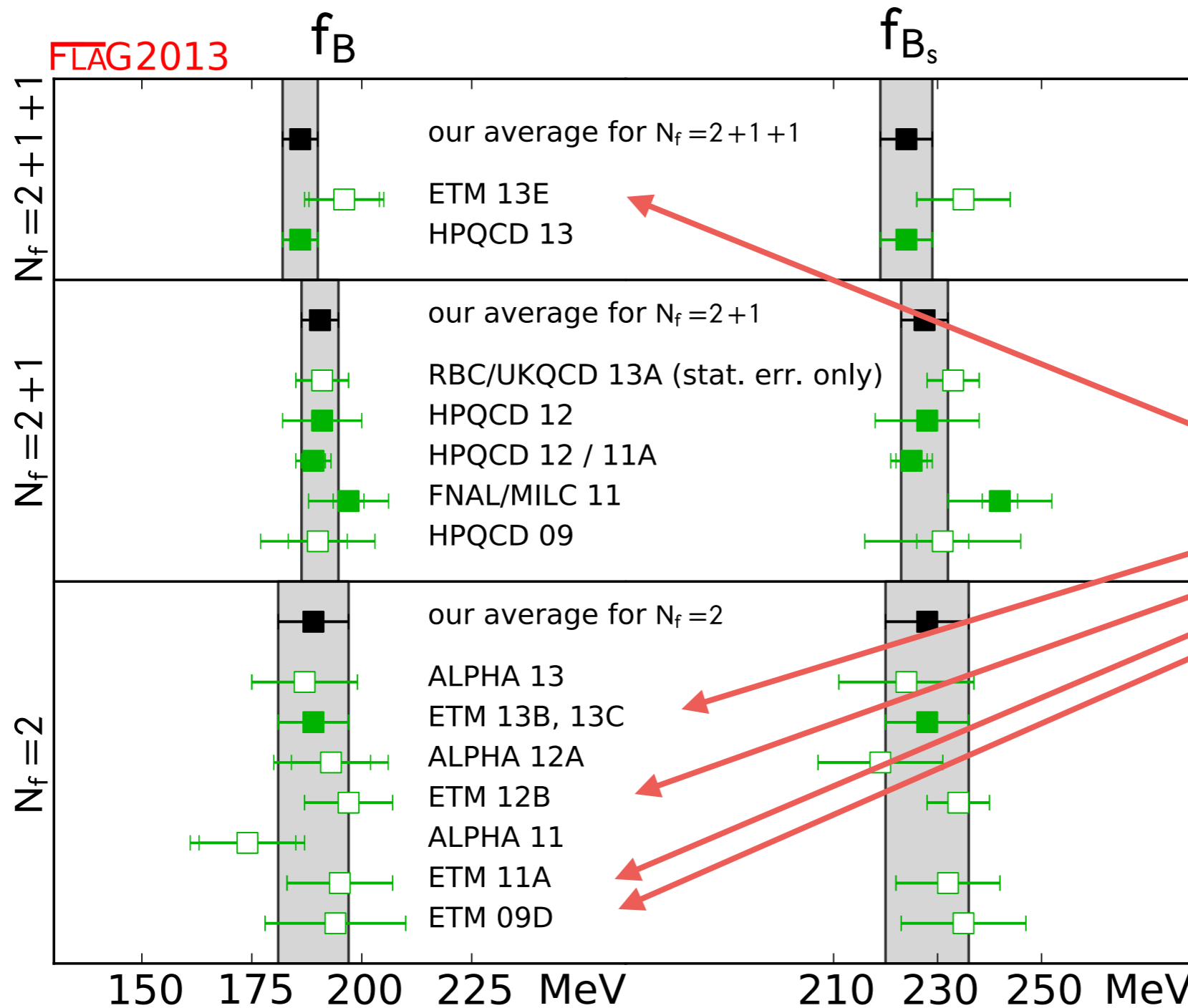
- final results

$$f_B = f_{hu/d}(\hat{\mu}_b) = 194(16)\text{MeV}$$

$$f_{B_s} = f_{hs}(\hat{\mu}_b) = 235(12)\text{MeV}$$

- error budget
 - L order $\sim 1\%$, NNL order $\sim 0.1\%$
 - ζ depends $\log x \sim 1, 2\%$
 - uncertainty of m_b

FLAG: f_B, f_{B_s}



ratio method

Summary

- Ratio method can deal with B-quark on current lattices.
- The procedure seems complicated, but good control of systematic errors.
- key factor - ratio, interpolation with static limit
- Is this method promising for B-physics ?

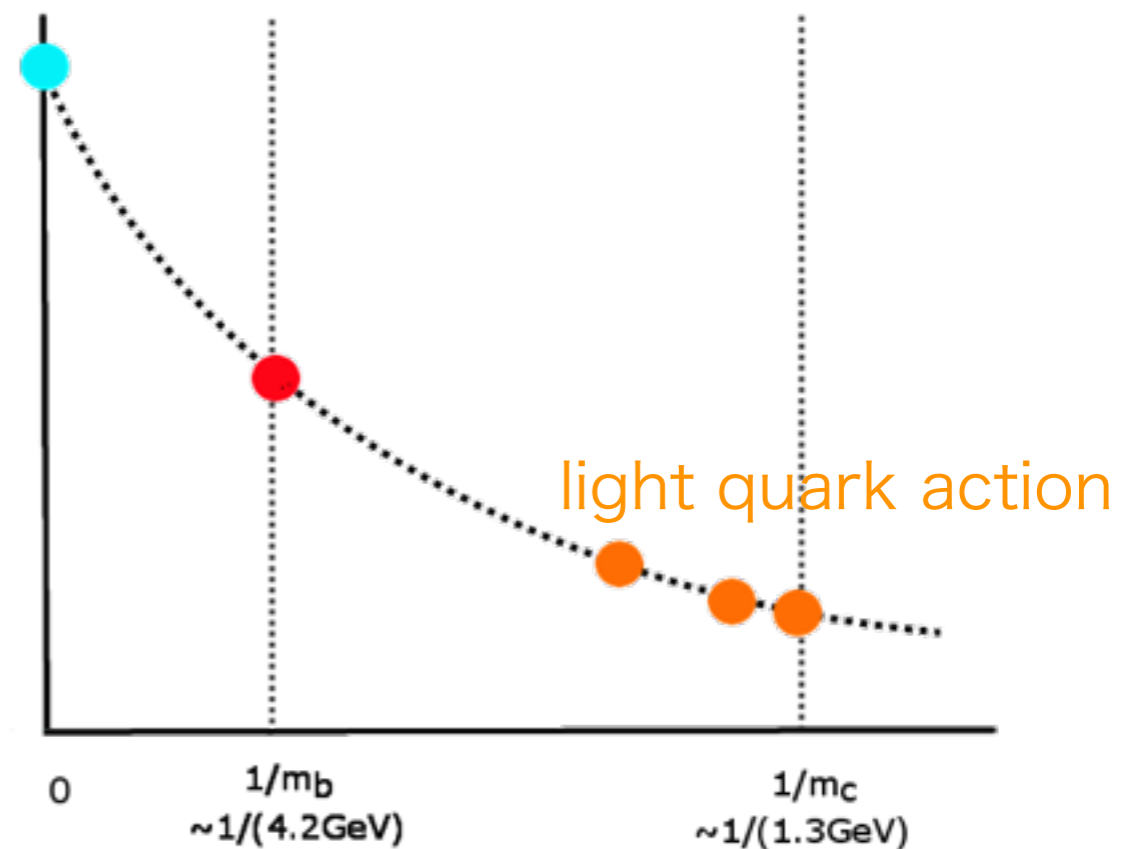
interpolation method

[JHEP01 (2012) 046]

- So far results of static limit is used.
- Lattice action of static limit is employed.

- average for
ratio method
interpolation method

lattice HQET



interpolation between lattice HQET and QCD data

$$S_{stat} = a^4 \sum_x \bar{\psi}_h(x) D_0 \psi_h(x) = a^3 \sum_x \bar{\psi}_h(x) [\psi_h(x) - V_{HYP}^\dagger(x - a\hat{0}, 0) \psi(x - a\hat{0})]$$

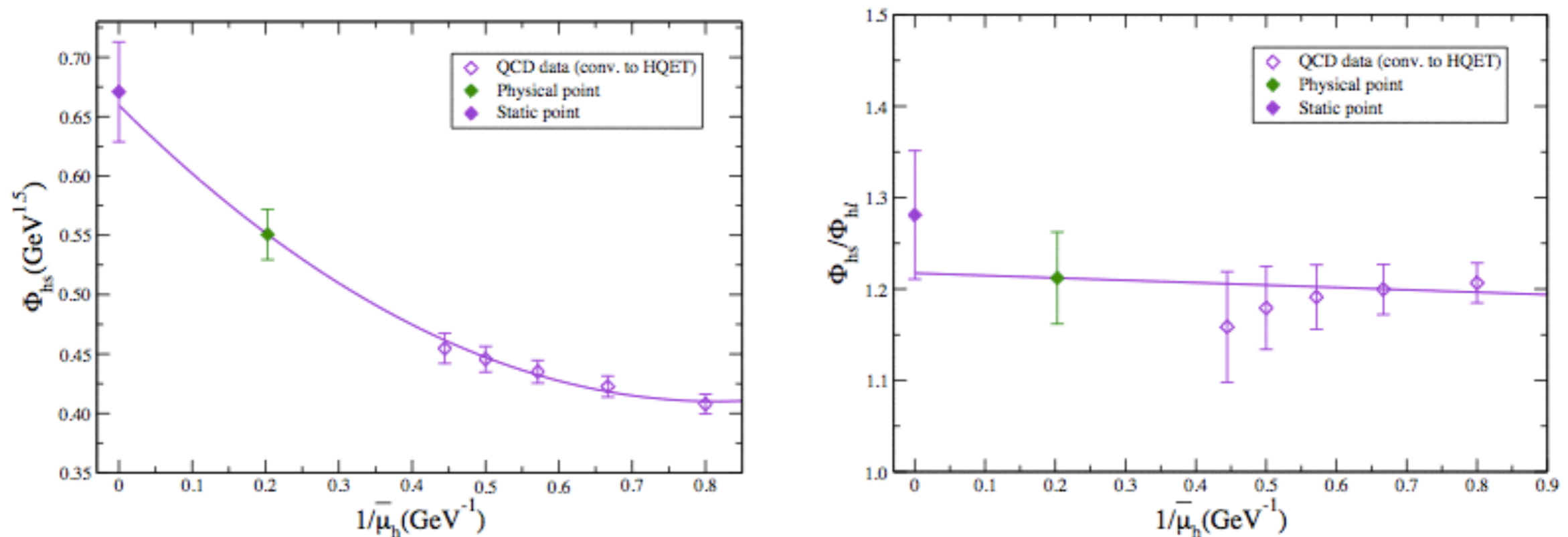


Figure 8: Dependence of Φ_{hs} (left) and Φ_{hs}/Φ_{hl} (right), in the chiral and continuum limit, on the inverse of the heavy quark mass.

final results

f_{B_s} [MeV]				f_{B_s}/f_B			
Ratio Method		Interpol. Method		Ratio Method		Interpol. Method	
Lin.	Quad.	Lin.	Quad.	HMChPT	Polyn.	HMChPT	Polyn.
225(7)(4)	225(7)(4)	237(9)(4)	238(9)(4)	1.22(2)(0)	1.14(2)(0)	1.22(5)(2)	1.16(6)(2)
225(7)(4)		238(9)(4)		1.18(2)(4)		1.19(5)(3)	
232(10)				1.19(5)			

Table 2: Collection of the results obtained for f_{B_s} and f_{B_s}/f_B from the ratio and interpolation methods. The statistical and systematic uncertainties are summed in quadrature. The third and fourth lines provide info on the results obtained by extrapolating to the physical pion mass point by using different chiral fit ansatz (see text). The final values, given in the last row, are an average of the results of the two methods.

HeavyMesonChPT

- SU(2) Chiral effective theories for heavy-light
- fixed heavy quark mass
- For pseudo scalar meson mass

$$M_{hl}r_0 = C_0 + C_1\hat{\mu}_l r_0 + \frac{a^2}{r_0^2}C_L$$

- For decay constant of PS meson

$$f_{hl}r_0 = D_0 + D_1\hat{\mu}_l r_0 + d_1 \frac{2B_0\hat{\mu}_l}{(4\pi f_0)^2} + \frac{a^2}{r_0^2}D_L$$

An improved ratio method for decay constant [JHEP03,016]

- Old method depends the measurement of the b-quark mass.
- propose a following asymptotic behavior.

$$\lim_{m \rightarrow \infty} \sqrt{M_{hl}} f_{hl} = \text{const}$$

- lattice ratio

$$z(x, \lambda; \mu) = \frac{f_{hl}(1/x)}{f_{hl}(1/\lambda x)} \cdot \frac{C_A^{stat}(\lambda x)}{C_A^{stat}(x)} \cdot \frac{M_{hl}(1/x)}{M_{hl}(1/\lambda x)}$$